



# Introduction to Nuclear and Particle Physics - 8.701 Lecture 2 S/W processes and Relativistic Kinematics

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## Lecture 2: Standard Model Interactions and Relativistic Kinematics

### Overview:

#### 1. Standard Model interactions:

- a.) Four forces
- b.) QED
- c.) QCD
- d.) Weak interactions

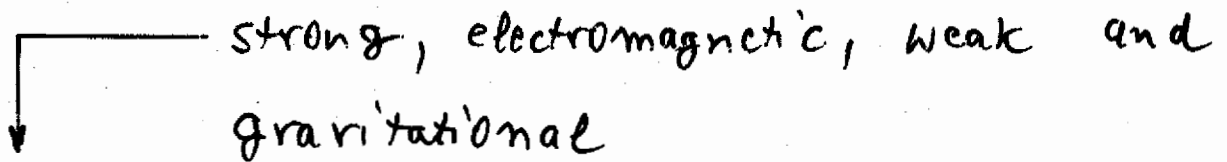
#### 2. Relativistic kinematics:

- a.) Lorentz transformation
- b.) Consequences of Lorentz transf.
- c.) Four vectors
- d.) Energy + momentum
- e.) Conservation laws

1. Standard Model Interactions:

a.) Four forces:

As far as we know, there are just four fundamental forces in nature:



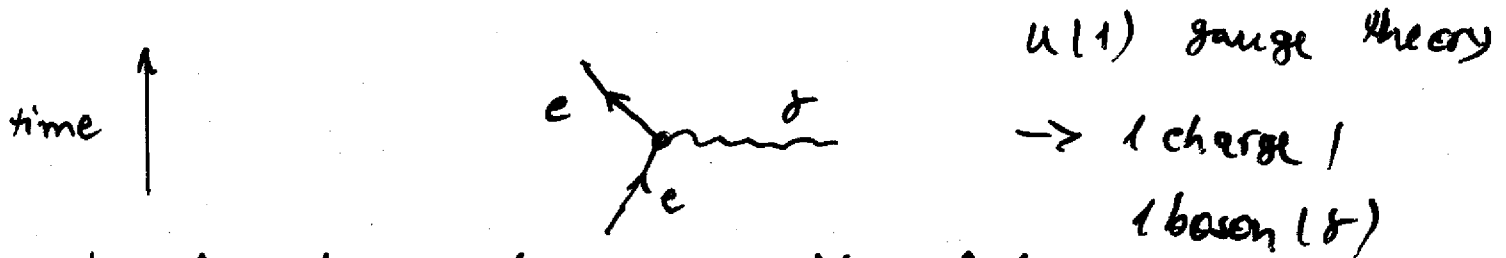
The Standard Model refers to three of these fundamental forces!

Properties:

	strong	electromag.	weak
Mediator	gluon ( $m=0$ )	$\gamma$ ( $m=0$ )	$W^{\pm}, Z^0$ <small><math>\sim 80\text{GeV}</math> <math>\sim 90\text{GeV}</math></small>
Acts on	color charge	el. charge	fermion
Particles experiencing it	gluons, quarks	charged part.	leptons, quarks
Range	$\lambda = \frac{h}{m\omega c} \sim 1\text{F}$	$\infty$	$\lambda = \frac{h}{m\omega c} \sim 10^{-3}\text{F}$
Typical life time	$10^{-23}\text{ s}$	$10^{-20} - 10^{-16}\text{ s}$	$10^{-12}\text{ s}$
Typical cross-section	$10\text{ mb}$ $\pi\text{p} \rightarrow \pi\text{p}$	$10^{-3}\text{ mb}$ $\sigma\text{p} \rightarrow \text{p}\pi^0$	$10^{-14}\text{ mb}$ $\nu\text{p} \rightarrow \nu\text{p}$
Typical coupling $\alpha_i$	1	$10^{-2}$	$10^{-6}$

6.) QED

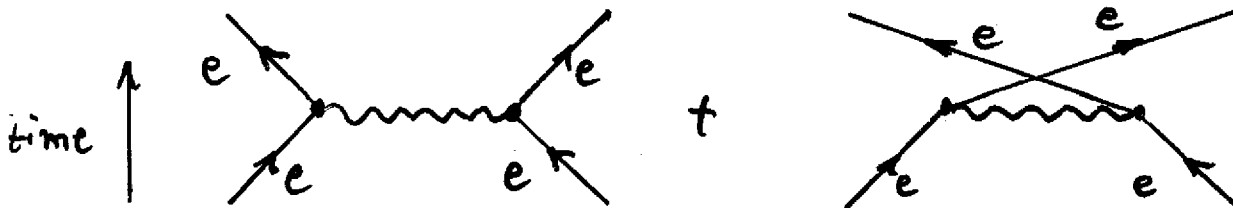
All electromagnetic phenomena are ultimately reducible to the following elementary coupling:



- two elements: charged particle; photon

- Meaning: charged particle  $e$  enters, emits (or absorbs) a photon  $\delta$  and exits

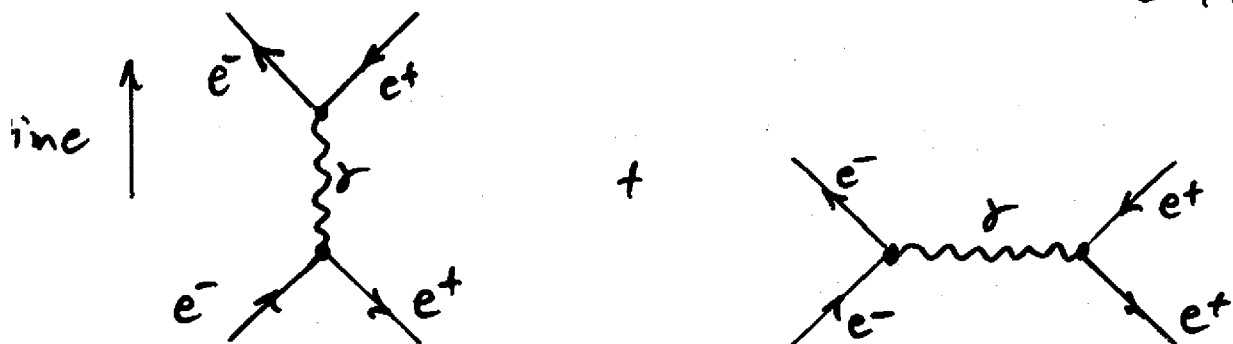
1. Example of a complete process: Møller scattering  
 $e + e \rightarrow e + e$



Meaning: interaction of two electrons which is mediated by a photon

classical case: Coulomb repulsion

2. Example of a complete process: Bhabha scattering  
 $e^- + e^+ \rightarrow e^- + e^+$



Rule: Particles which are running "backward in time" are to be interpreted as the corresponding anti-particle!

Bhabha scattering is related to mueller scattering by a general principle which is known as crossing symmetry:

$$A + B \rightarrow C + D$$

Rule: Any particle can be "crossed" over the other side of the equation, provided it turns into its anti-particle:

$$A \rightarrow \bar{B} + C + D$$

$$A + \bar{C} \rightarrow \bar{B} + D$$

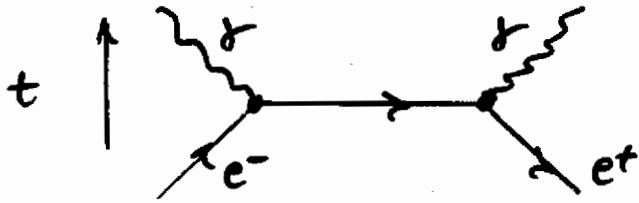
$$\bar{C} + \bar{D} \rightarrow \bar{A} + \bar{B}$$

} These processes are dynamically allowed, but not necessarily kinematically!

(e.g. if A weighs less than  $\bar{B}$ , C and D, then this process is kinematically not allowed)

Bhabha scattering and mueller scattering are related by crossing symmetry!

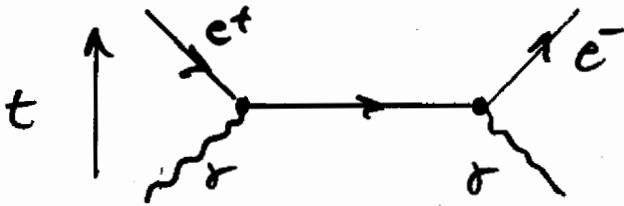
3. Example of complete process:



Pair annihilation

$$e^- + e^+ \rightarrow \gamma + \gamma$$

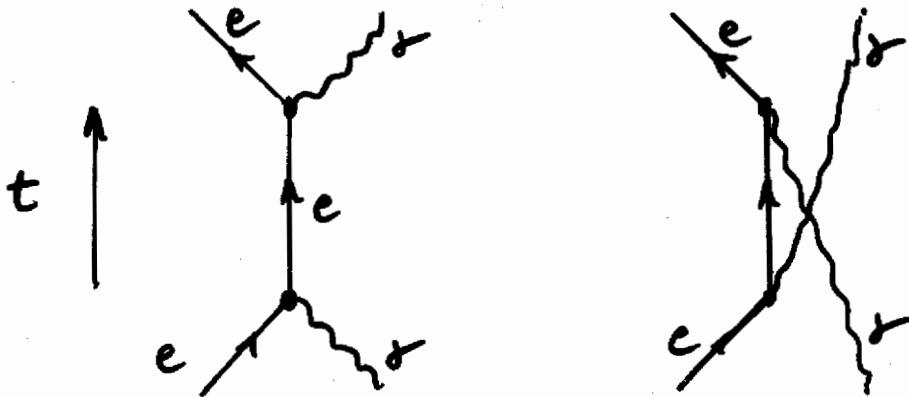
4. Example of complete process:



Pair production

$$\gamma + \gamma \rightarrow e^- + e^+$$

5. Example of complete process:

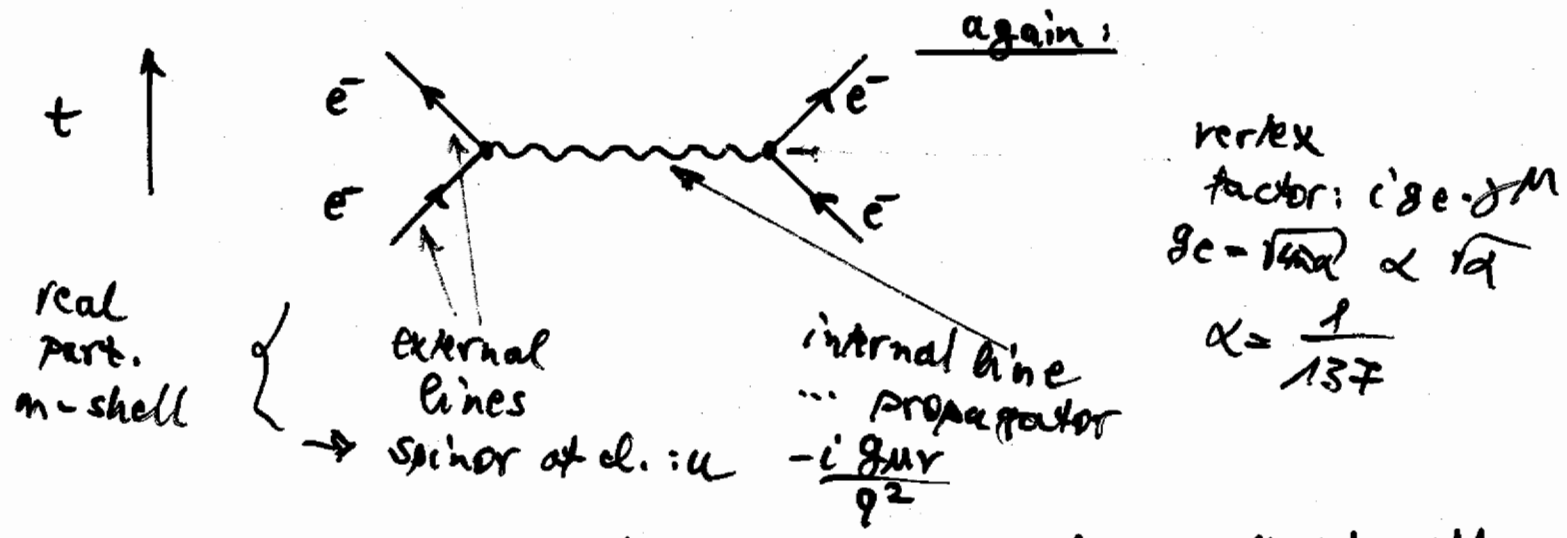


Compton scattering

$$e + \gamma \rightarrow e + \gamma$$

Remarks on Feynman diagrams:

1. Feynman diagrams are purely symbolic, they do not represent particle trajectories.
2. Example: "Components" of a Feynman diagram

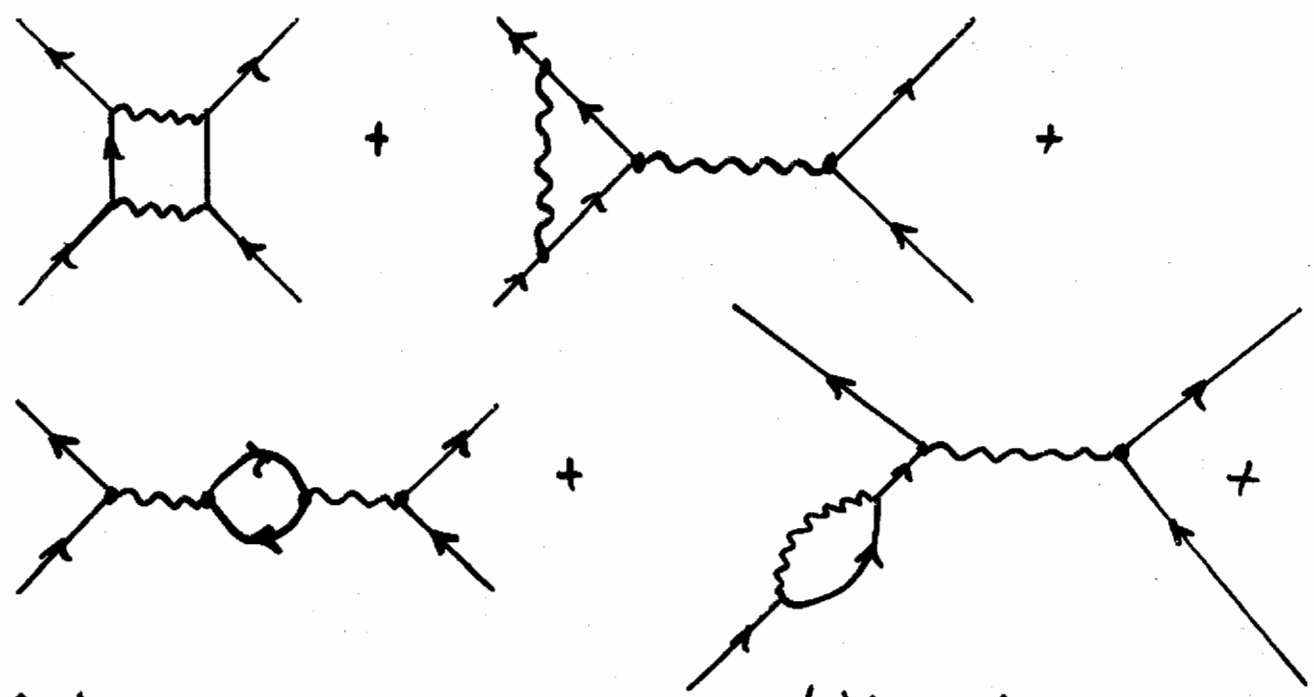


3. Each Feynman diagram represents a Matrix,  $U$ , to account for a particular process
4. Feynman rules enforce conservation of energy and momentum at each vertex, and hence for the diagram as a whole.
5. Procedure:
  - a) Draw all diagrams that have the appropriate external lines
  - b) Sum of all diagrams with the given external lines represents the actual physical process.

Problem: There are infinite many Feynman diagrams!

Example: Møller scattering

Besides the above 2-vertex Feynman diagrams for Møller scattering, we have with 4 vertices:



calculations get very complicated!

higher order diagrams

BUT: Each vertex contributes a factor  $\alpha$ .

$\alpha$ : QED Feynman structure constant (general: QED coupling constant)

Higher order terms are suppressed since  $\alpha \ll 1$  (Perturbation theory)!

What matters "mainly" is the leading order contribution!

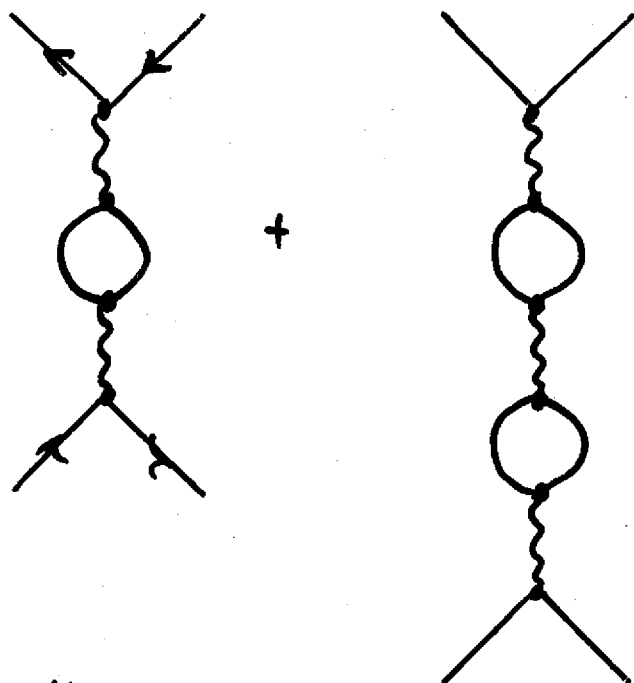
However: There are many examples where measurements are so precise that higher order terms have to be taken into account!  $\rightarrow$  TEST OF QED

example: Measurement of anom. magn. moment of muon!

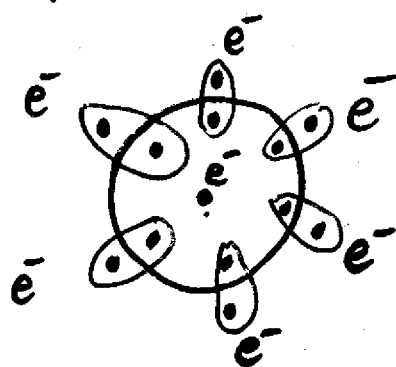
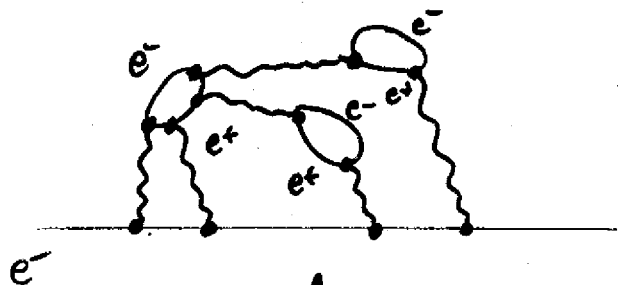


QED vacuum polarization:

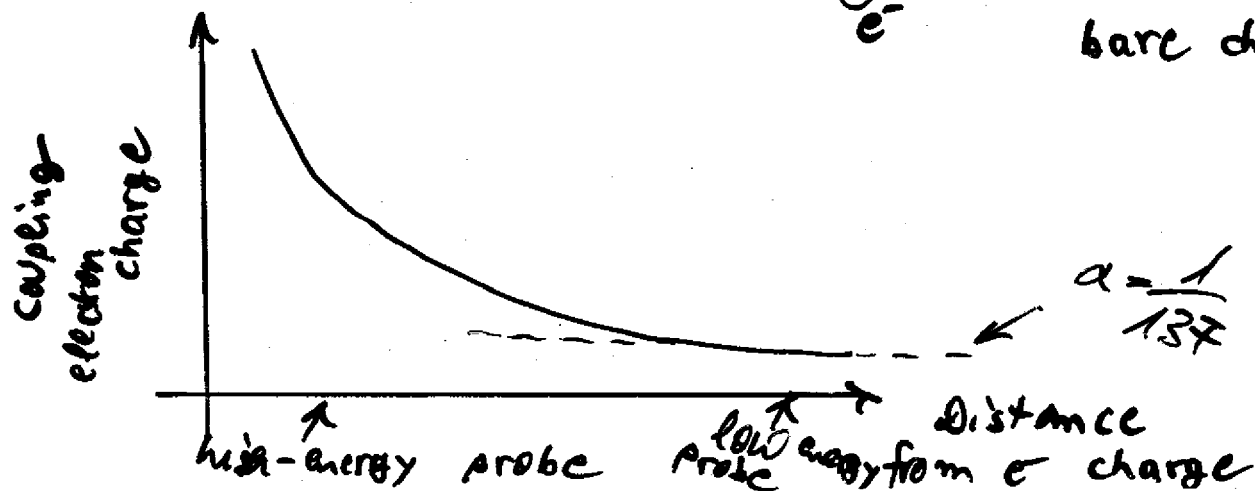
In QED, the vacuum behaves like a dielectric, creating  $e^-/e^+$  pairs out of the vacuum:



or view it like this:



"bare" electron charge is screened  
eff. charge  $\ll$  bare charge



c.) QCD :

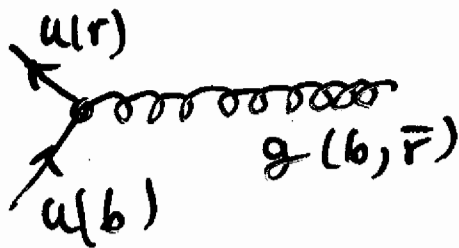
In Quantum - Chromo - Dynamics :  $SU(3)$  gauge theory

→ 3 charges / 8 bosons (gluons)

color plays the role of charge and the fundamental

process is: quark → quark + gluon

• Fundamental vertex :



color is always

conserved:

→ gluon is carrying away the difference

• Types of gluons: 8 "possibilities"

$r\bar{r}, r\bar{b}, r\bar{g}, b\bar{r}, b\bar{b}, b\bar{g}, g\bar{r}, g\bar{b}, g\bar{g}$

color octet  $11\gamma = (1/\sqrt{2})(r\bar{b} + b\bar{r})$

$15\gamma = -i(r\bar{g} - g\bar{r})/\sqrt{2}$

color octet  $12\gamma = -1/\sqrt{2}(r\bar{b} - b\bar{r})$

$16\gamma = (b\bar{g} + g\bar{b})/\sqrt{2}$

color octet  $13\gamma = (r\bar{r} - b\bar{b})/\sqrt{2}$

$17\gamma = -i(b\bar{g} - g\bar{b})/\sqrt{2}$

color octet  $14\gamma = (r\bar{g} + g\bar{r})/\sqrt{2}$

$18\gamma = (r\bar{r} + b\bar{b} - 2g\bar{g})/\sqrt{6}$

color singlet  $19\gamma = (r\bar{r} + b\bar{b} + g\bar{g})/\sqrt{3}$

"no net color"

... like a photon...

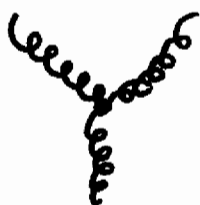
Note:

Confinement requires that all naturally occurring particles be color singlets.  $197$  is a color singlet state. If it exists as a mediator, it should also occur as a free particle  
→ As such a mediator: Exchange between color singlet particles, e.g. a proton and neutron  $\Rightarrow$  long-range force with strong coupling.

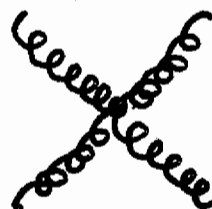
However: strong force is of very short-range. So: Experiments tell us that there are only 8 gluons: color octet:  $SU(3)$

Because gluons themselves carry color (in contrast to the photon which is electrically neutral): gluons couple directly to one another:

3 gluon vertices:



4 gluon vertices:



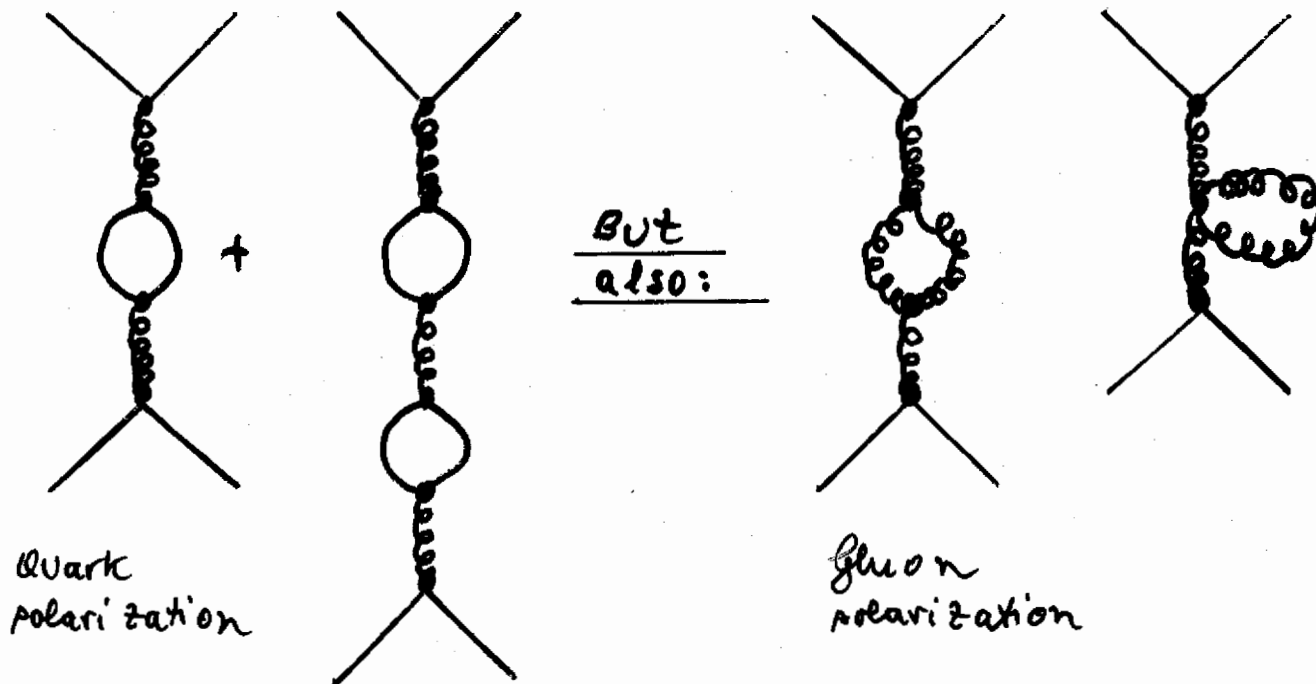
Calculations: Apply QCD Feynman rules to calculate various processes!

→ How does the QCD coupling constant behave compared to the QED coupling constant?

Clear difference between QED and QCD:

→ Self-coupling of gluons which is absent in case of QED.

QCD vacuum polarization diagrams:



→ Behavior is opposite for  $\alpha_s$  :

Quark polarization:  $\alpha_s$  is large at short distances

gluon polarization:  $\alpha_s$  is small at short distances

A priori not clear, who wins!

The Winner depends on the relative number of flavors ( $f$ ) and colors ( $n$ ):

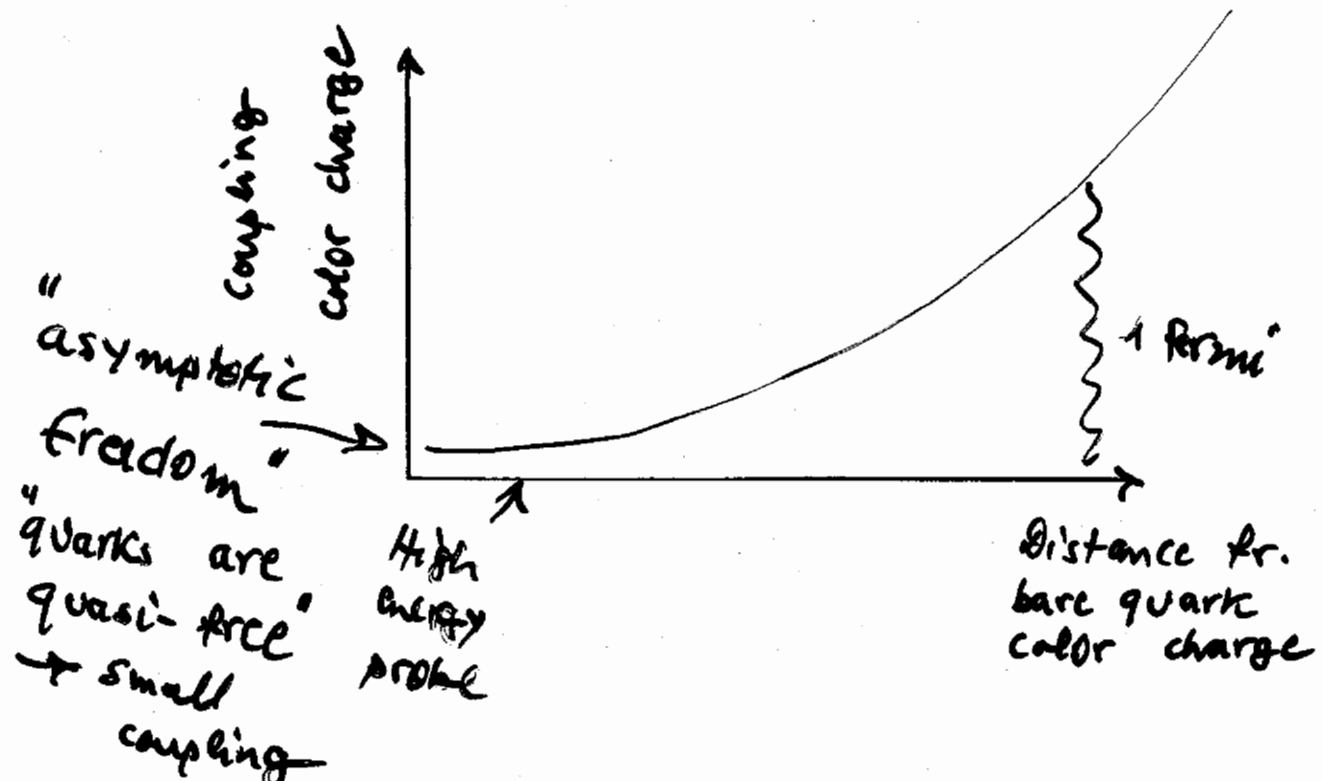
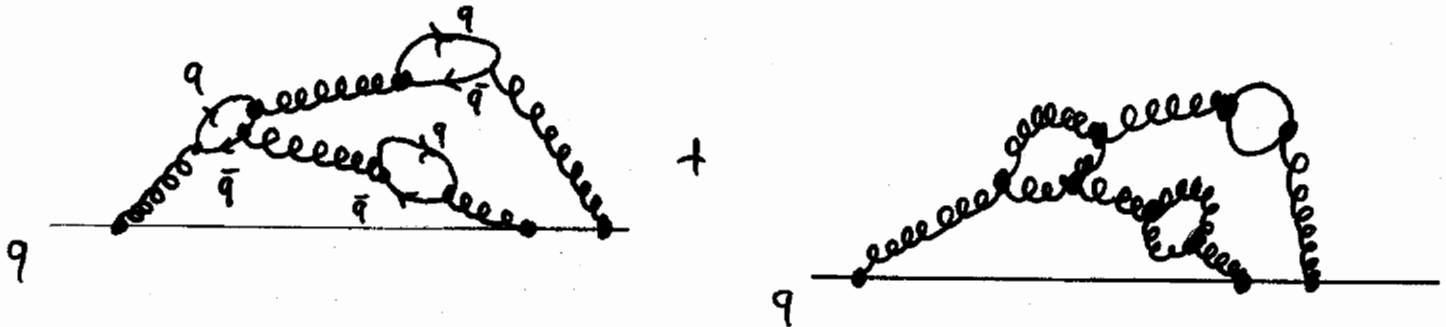
critical parameter:

$$a \equiv 2f - 11n$$

Standard Model:  $f = 6, n = 3 \rightarrow a = -21$

→ QCD coupling constant decreases at short distances!

Another "last picture":



d) Weak interactions:

SU(2) gauge theory: 3 mediators:  $Z^0$ ;  $W^\pm$

- Quarks and leptons take part in weak interactions

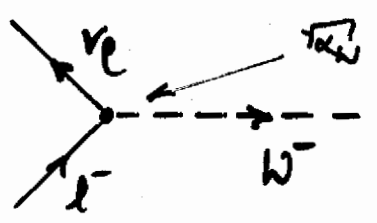
- 2 types of interactions:

a.) charged current:  $W^\pm$

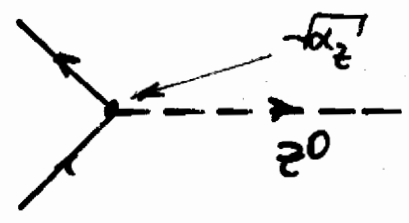
b.) neutral current:  $Z^0$

1. leptons:

charged vertex:



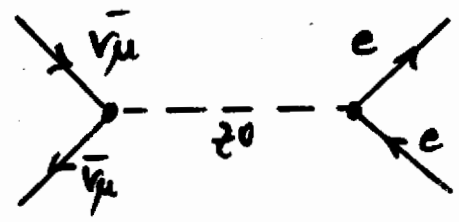
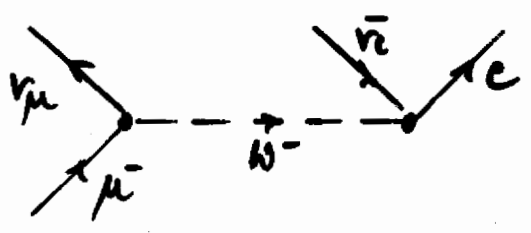
neutral vertex:



example:

$$\mu^- \rightarrow e^- + \nu_\mu + \bar{\nu}_e$$

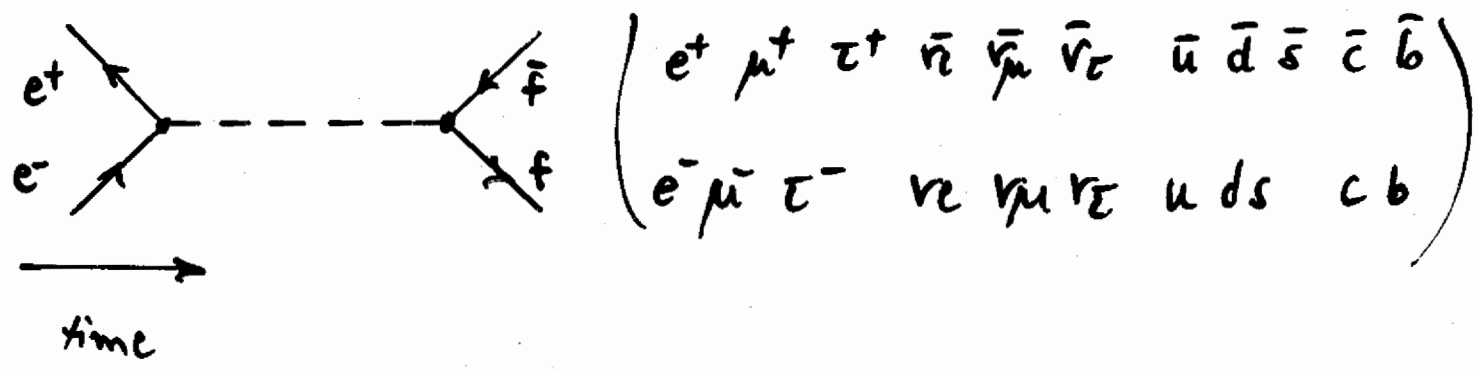
$$\bar{\nu}_\mu + e^- \rightarrow \bar{\nu}_\mu + e^-$$



First "picture" of neutral weak process discovered at CERN in 1973.

The Glashow-Weinberg-Salam (GWS) model includes neutral weak processes as an essential ingredient. Their existence was confirmed experimentally at CERN in 1973.

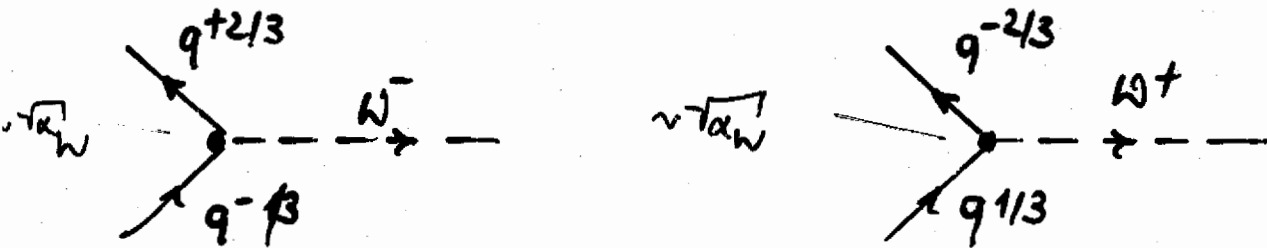
Production of  $Z^0$  at LEP : max.  $E_{e^-} + E_{e^+}$  at  $m_{Z^0}$



established 3 families in nature!

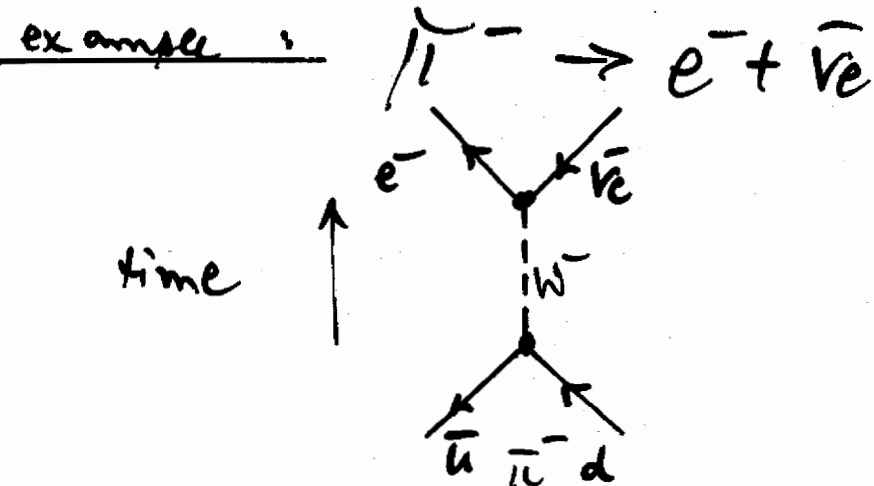
2. Quarks:

Fundamental vertex:  $W^\pm$



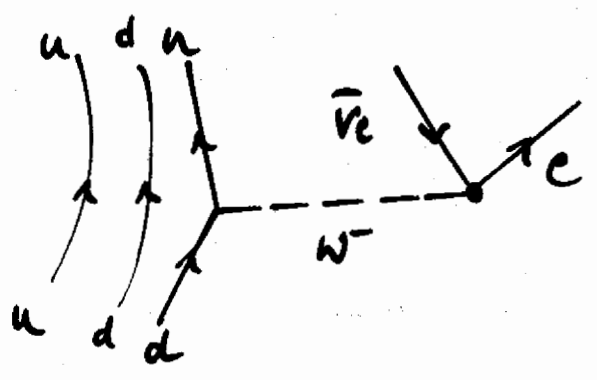
note: A quark of charge  $-1/3$  ( $d, s, b$ ) converts into the corresponding quark with charge  $+2/3$  ( $u, c, t$ ) with the emission of  $W^-$  (vice versa for  $W^+$ ).

Flavor is not conserved in weak interactions! Since the quark flavor changes at a weak vertex ( $W^\pm$ ), as a quark color changes at a strong vertex, weak interactions are sometimes called "flavor dyn."



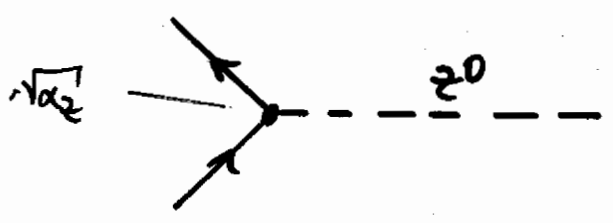


Beta-decay of the neutron:



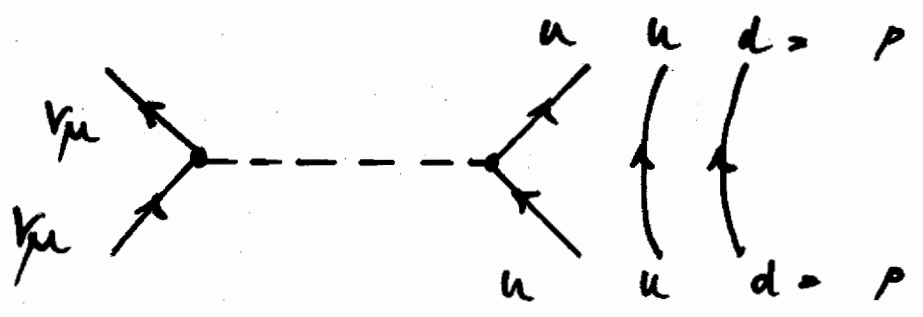
$$n \rightarrow p + e^- + \bar{\nu}_e$$

Fundamental vertex: Z0



note: quark flavor is not changed!

Example:



$$\nu_\mu + \bar{\nu}_\mu \rightarrow \nu_\mu + p$$

o Flavor changing reactions:  $W^\pm$

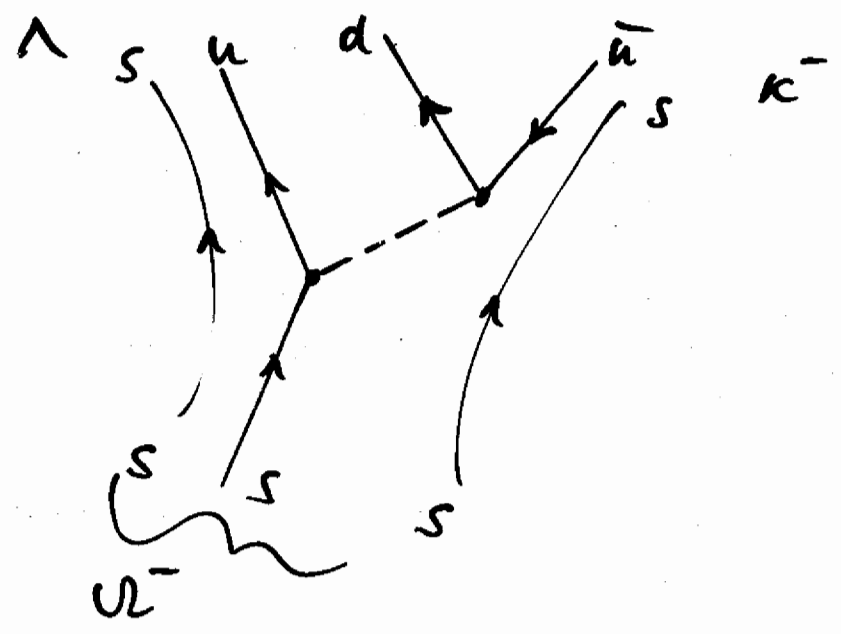
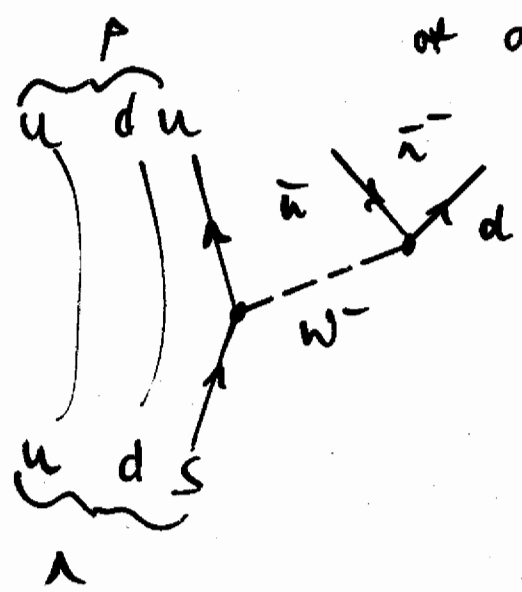
In the spirit of the charged weak coupling with respect to leptons which yield only changes within each family, i.e. :  $e^- \leftrightarrow \nu_e$  ;  $\mu^- \leftrightarrow \nu_\mu$  ;  $\tau^- \leftrightarrow \nu_\tau$

one might assume that this also holds for quarks:

$$\begin{pmatrix} u \\ d \end{pmatrix} \quad \begin{pmatrix} c \\ s \end{pmatrix} \quad \begin{pmatrix} t \\ b \end{pmatrix}$$

However:

The observed decay  $\Lambda \rightarrow p + \bar{u}^+$  or  $\bar{K}^0 \rightarrow \Lambda + K^-$  involve the conversion of a strange quark into an up-quark:



note: Flavor changes do occur not only within one family!

Therefore:

The flavor eigenstate  $|u\rangle$  is not the partner to the flavor eigenstate  $|d\rangle$ , but to a linear combination of  $d$ ,  $s$  and  $b$ :

$$\begin{pmatrix} u \\ d' \end{pmatrix} = \begin{pmatrix} c \\ s' \end{pmatrix} = \begin{pmatrix} t \\ b' \end{pmatrix}$$

$|d'\rangle$  can be expressed as a linear combination of  $|d\rangle$ ,  $|s\rangle$  and  $|b\rangle$ :

$$\begin{pmatrix} |d'\rangle \\ |s'\rangle \\ |b'\rangle \end{pmatrix} = \begin{pmatrix} \underline{V_{ud}} & \underline{V_{us}} & \underline{V_{ub}} \\ \underline{V_{cd}} & \underline{V_{cs}} & \underline{V_{cb}} \\ \underline{V_{td}} & \underline{V_{ts}} & \underline{V_{tb}} \end{pmatrix} \begin{pmatrix} |d\rangle \\ |s\rangle \\ |b\rangle \end{pmatrix}$$

Note:

The matrix is called after their "inventors":

3 x 3 Kobayashi - Maskawa Matrix! The coefficients are sometimes expressed as cosine and sine values of an angle: Cabibbo angle  $\theta_c$

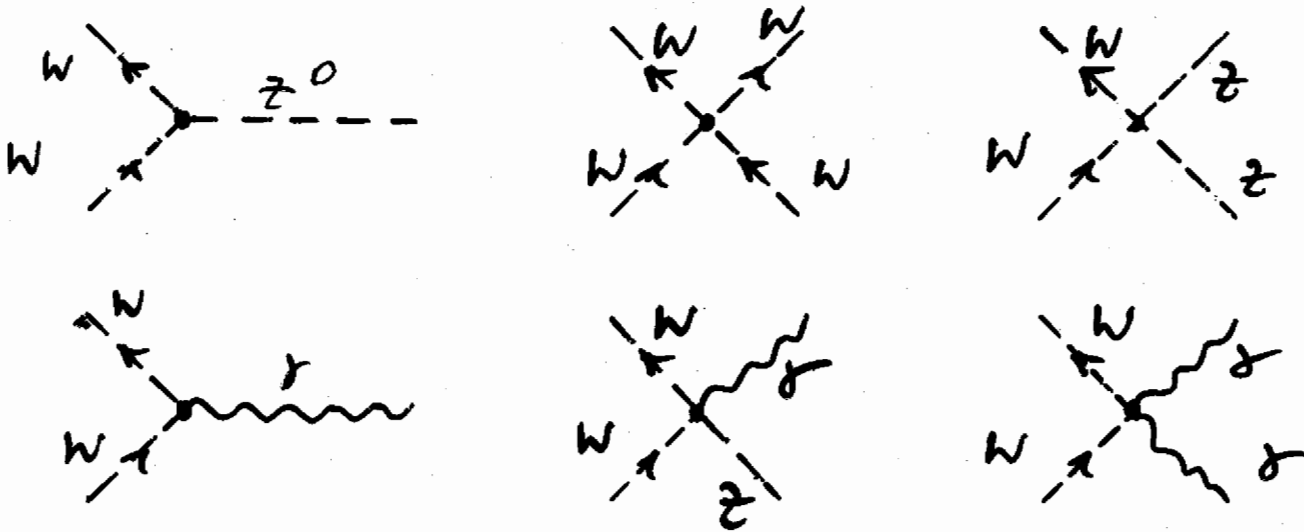
earlier scheme (2 families):  $|d'\rangle = |d\rangle \cos \theta_c + |s\rangle \sin \theta_c$   
 $|s'\rangle = |d\rangle \sin \theta_c + |s\rangle \cos \theta_c$

Experimental data:

$$V_{ij} = \begin{pmatrix} 0.9741 & 0.0219 & 0.0005 \\ 0.0256 & -0.0226 & -0.0014 \\ 0.019 & 0.0372 & 0.0038 \\ 0.021 & -0.0349 & -0.0014 \\ 0.004 & 0.037 & 0.0036 \\ 0.014 & -0.044 & -0.0017 \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

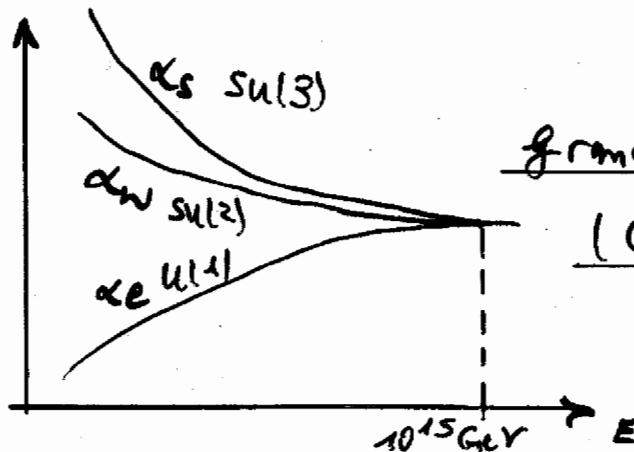
K. Hagiwara et al., Phys. Rev. D 66 (2002) 010001.

Self coupling of weak-bosons and photon:



Last remark:

$\alpha_i$   
 $su(3) \otimes su(2) \otimes u(1)$

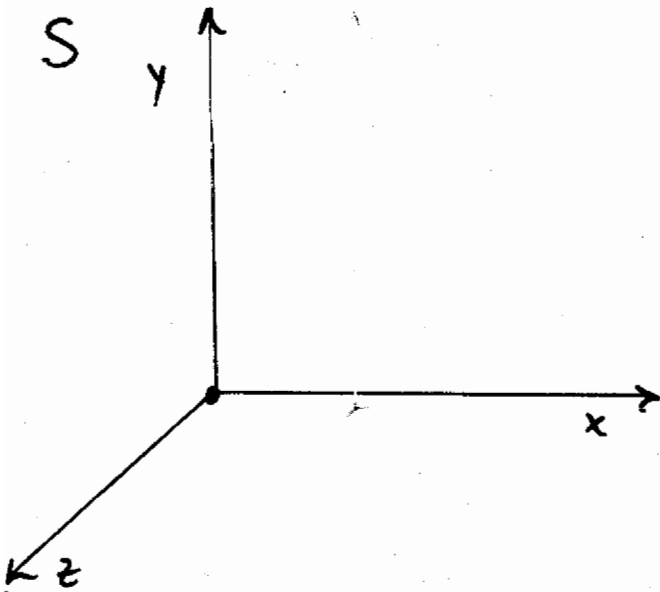


Grand Unified Theories  
 (GUT) "one Force"

## 2. Relativistic kinematics:

### a.) Lorentz transformations:

Given two inertial frames  $S$  and  $S'$ , with  $S'$  moving at uniform speed  $v$  with respect to  $S$ :



Events are related in space-time as follows

(Lorentz transformations):

$$x' = \gamma(x - vt)$$

$$x = \gamma(x' + vt')$$

$$y' = y$$

$$y = y'$$

$$z' = z$$

$$z = z'$$

$$t = \gamma\left(t - \frac{v}{c^2} \cdot x\right)$$

$$t = \gamma\left(t' + \frac{v}{c^2} \cdot x'\right)$$

With:  $\gamma = 1 / \sqrt{1 - v^2/c^2}$

$$\beta = (v/c)$$

1.) Consequences:

1. relativity of simultaneity:

If two events occur at the same time in  $S$ , but at different locations, then they do not occur at the same time in  $S'$ : With  $t_A = t_B$ :

$$t'_A = t'_B + \frac{\gamma v}{c^2} (x_B - x_A)$$

2. Lorentz contraction:

A moving object is shortened by a factor  $\gamma$ :

$$L = L' \cdot \gamma^{-1}$$

3. Time dilation:

Moving clocks run slow:

$$T = T' \cdot \gamma$$

4. velocity addition:

Particle moves with speed  $u$  with respect to  $S'$ .  
What is the speed,  $u$ , with respect to  $S$ ? Note:  
 $S'$  moves with speed  $v$  with respect to  $S$ :

$$u = \frac{u' + v}{1 + (u' \cdot v / c^2)}$$

if  $u' = c \rightarrow u = c$

c.) Four vectors:

$$x^0 = c \cdot t ; x^1 = x ; x^2 = y ; x^3 = z$$

Transformation:

$$x^{0'} = \gamma(x^0 - \beta x^1)$$

$$x^{1'} = \gamma(x^1 - \beta x^0)$$

$$x^{2'} = x^2$$

$$x^{3'} = x^3$$

$$\beta = v/c$$

compact:  $x^{\mu'} = \sum_{\nu=0}^3 \Lambda^{\mu}_{\nu} \cdot x^{\nu}$   
 $\mu = 0, 1, 2, 3$

$$\Lambda = \begin{pmatrix} \gamma & -\gamma\beta & 0 & 0 \\ -\gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

short:

$$x^\mu = \Lambda^\mu_\nu \cdot x^\nu$$

convention by Einstein:

Repeated indices are to be summed!

Invariant:

$$I \equiv (x^0)^2 - (x^1)^2 - (x^2)^2 - (x^3)^2 = (x^{0'})^2 - (x^{1'})^2 + (x^{2'})^2 + (x^{3'})^2$$

→ Same in any inertial system!

Different notation:

$$g = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$

$$I = g_{\mu\nu} \cdot x^\mu \cdot x^\nu$$

$g_{\mu\nu}$ : Metric tensor

Covariant four vector:

$$x_\mu = g_{\mu\nu} \cdot x^\nu$$

Contravariant four vector:

$$x^\mu \quad (I = x_\mu x^\mu)$$



With any two four vectors:

$$a^\mu \cdot b_\mu = a_\mu \cdot b^\mu = a^0 b^0 - a^1 b^1 - a^2 b^2 - a^3 b^3$$

W/  $b_\mu = a_\mu$  :

$$a^2 \equiv a \cdot a = (a^0)^2 - (\vec{a})^2$$

$a^2 > 0$ : timelike

$a^2 < 0$ : spacelike

$a^2 = 0$ : lightlike

$$p_\mu \cdot p^\mu = \frac{E^2}{c^2} - (\vec{p})^2 = m^2 c^2$$

d.) Energy + momentum:

$$\vec{\eta} = \gamma \cdot \vec{v} \quad (\text{proper velocity}) : \eta^\mu = \gamma (c, v_x, v_y, v_z)$$

$$\vec{p} = m \cdot \vec{\eta} = \gamma \cdot m \vec{v} ; p^0 = \gamma m c$$

$$E = \gamma m c^2 = \frac{m c^2}{(1 - \beta^2)^{1/2}}$$

"relativistic energy"

$$p^\mu = \left( \frac{E}{c}, p_x, p_y, p_z \right)$$

e.) Conservation laws : analysis of particle

reactions :

Conservation of :

1. Energy and momentum

2. Angular momentum

3. Electric charge

4. Color charge

5. Baryon number :  $A=1$  for Baryons ;  $A=-1$  for Anti-Baryons  $A=0$  for non Baryons

6. Lepton number :

Particles at each generation ( leptons) are conserved :

$L_e, L_\mu, L_\tau$

$L_i = +1$  : particle

$L_i = -1$  : anti-particle

example :  $\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$

7. Flavor is conserved in strong and electro mag. interactions, but not in weak interactions ( $W^\pm$ ) !