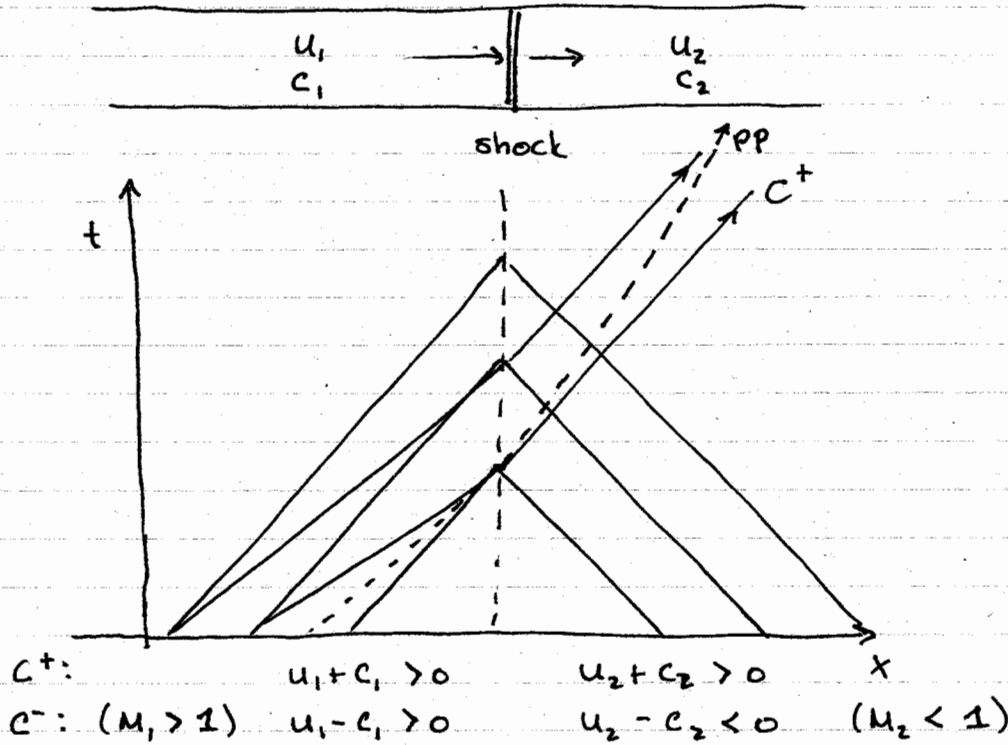


Lecture 10

Pset 3: 8.2

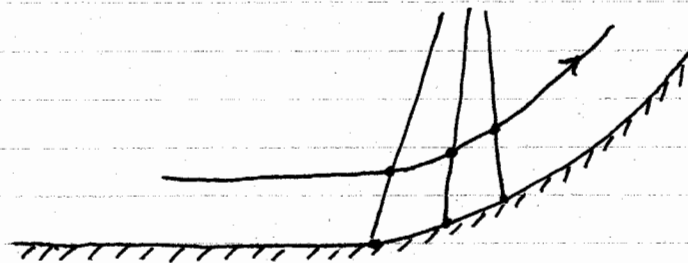
Normal shock (standing) in 1D flow



So far we've seen: quasi-1D steady
quasi-1D unsteady

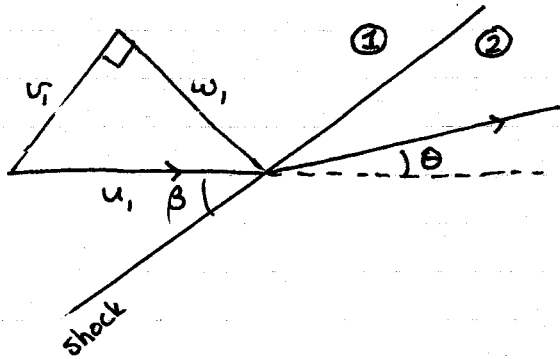
Today: 2D steady (chpt. 9)

Supersonic stream turns a corner through a series of standing oblique shocks:



↑ # of shocks
∞ # of weak shocks
⇒ isentropic turn

Recall for weak oblique shocks



$$[\theta] = - \frac{\sqrt{M_1^2 - 1}}{M_1^2} \frac{[\omega]}{c_1}$$

(1st term in Taylor exp.)

Since $v_1 = v_2$,

$$u_2^2 - u_1^2 = (\omega_2^2 + v_2^2) - (\omega_1^2 + v_1^2) = \omega_2^2 - \omega_1^2$$

$$(u_2 + u_1)[u] = (\omega_2 + \omega_1)[\omega]$$

$$(2u_1 + \overset{\text{small}}{[u]})[u] = (2\omega_1 + \overset{\text{small}}{[\omega]})[\omega] \Rightarrow u_1[u] \approx \omega_1[\omega]$$

From diagram: $\frac{\omega_1}{u_1} = \sin \beta$

Recall for weak oblique shocks $\beta \approx \mu_1 = \frac{1}{M_1}$

Mach angle

$$\frac{u_1[u]}{u_1} = \frac{\overset{1/M_1}{\omega_1}}{u_1} [\omega] \Rightarrow \frac{[u]}{u_1} = \frac{1}{M_1^2} \frac{[\omega]}{c_1}$$

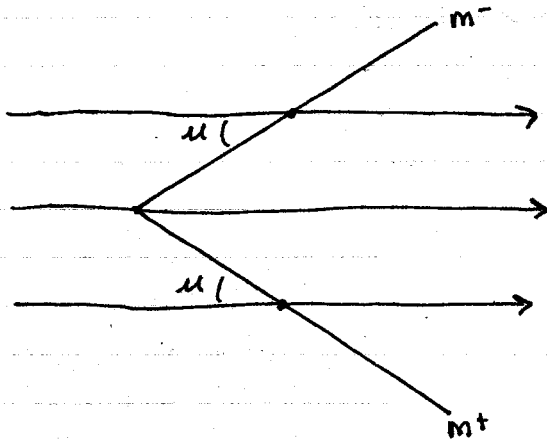
$$\Rightarrow [\theta] = -\sqrt{M_1^2 - 1} \frac{[u]}{u_1}$$

Infinitesimal strength shock:

$$d\theta = \pm \sqrt{M_1^2 - 1} \frac{du}{u_1}$$

↑

see diagram...



$$d\theta = -\sqrt{M^2-1} \frac{du}{u}$$

↷ θ

$$d\theta = +\sqrt{M^2-1} \frac{du}{u}$$

Recall for steady flow:

$$\frac{du}{u} = \frac{dM/M}{1+(\gamma-1)M^2}$$

define $dw = \pm d\theta$

$$dw = \frac{\sqrt{M^2-1}}{1+(\gamma-1)M^2} \frac{dM}{M}$$

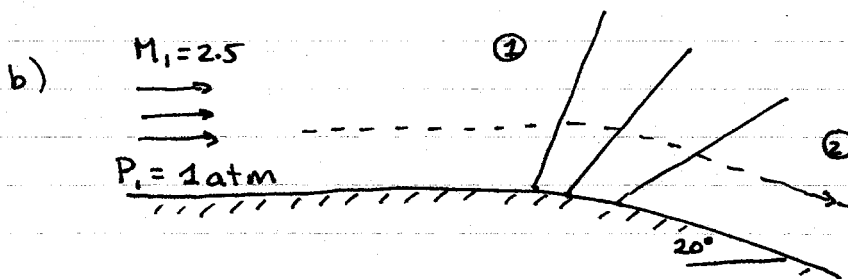
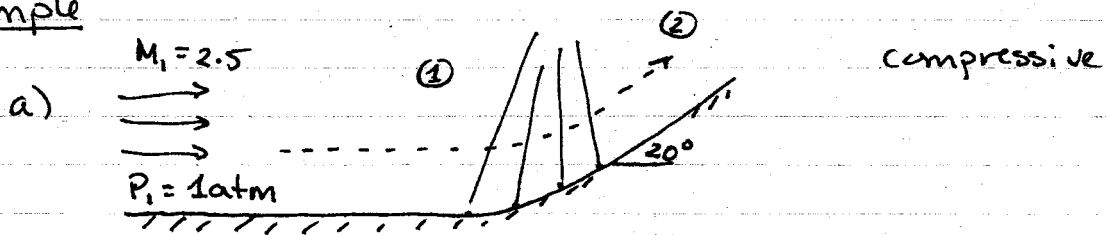
$w =$ Prandtl-Meyer Func.

For a perfect gas, this can be integrated:

$$w(M) = \sqrt{\frac{\gamma+1}{\gamma-1}} \tan^{-1} \sqrt{\frac{\gamma-1}{\gamma+1} (M^2-1)} - \tan^{-1} \sqrt{M^2-1}$$

Why is this useful? w is another dimensionless measure of flow speed. (alternative to the Mach #).

Example



a) $\Delta\omega = -\Delta\theta = -20^\circ$

$M = 2.5 \Rightarrow \omega = 39.12^\circ$

$\Rightarrow \omega_2 = 19.12^\circ \Rightarrow M_2 = 1.75$

b) $\Delta\omega = \Delta\theta = 20^\circ$

$M = 2.5 \Rightarrow \omega = 39.12^\circ$

$\Rightarrow \omega_2 = 59.12^\circ \Rightarrow M_2 = 3.54$

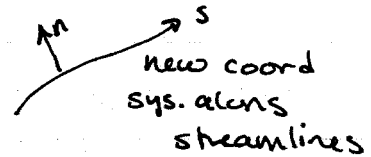
Method of Characteristics

Equations of motion:

$\nabla \cdot (\rho \vec{u}) = 0$ continuity

$\nabla \times \vec{u} = 0$ irrotationality

$\nabla \left(\frac{u^2}{2} \right) + \frac{1}{\rho} \nabla P = 0$ momentum



↓ lots of algebra

$\frac{d}{dm} (\theta + \omega) = \frac{\sin \mu \sin \theta}{r}$

$r \rightarrow \infty \Rightarrow$ plane flow

$\frac{d}{dm} (\theta - \omega) = - \frac{\sin \mu \sin \theta}{r}$

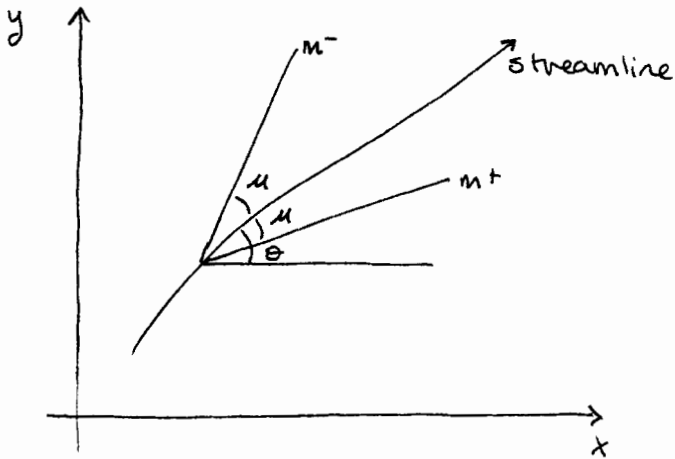
↑ ↑
u thermodynamics

•
↙ ↘
along streamline

↙ ↘
normal to streamline

where: $\frac{d}{dm} = \cos \mu \left(\frac{\partial}{\partial s} + \tan \mu \frac{\partial}{\partial n} \right)$

$\frac{d}{dm} = \cos \mu \left(\frac{\partial}{\partial s} - \tan \mu \frac{\partial}{\partial n} \right)$

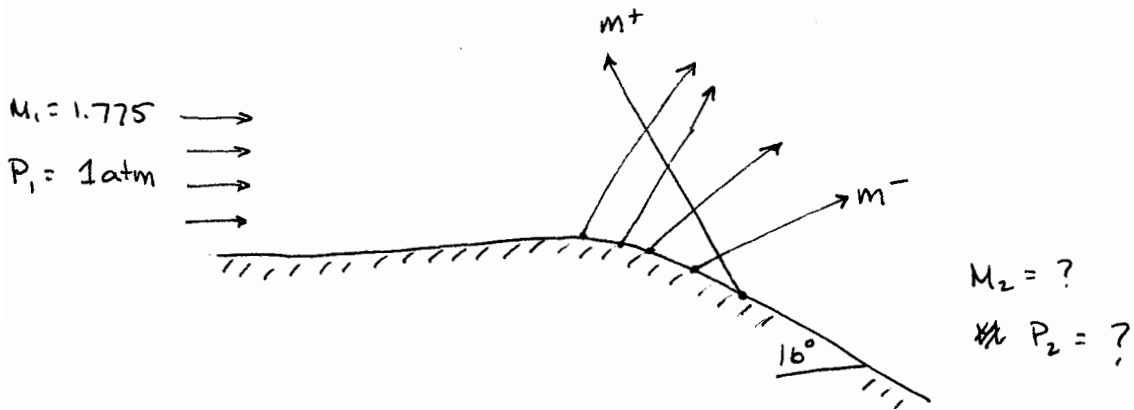


(note axes are y-x
not t-x. analogous
 but not identical to
 1D case.)

Invariants:

$$\left. \begin{aligned} \theta + \omega &= \text{const. along } m^+ \\ \theta - \omega &= \text{const. along } m^- \end{aligned} \right\} \text{ for planar flow}$$

Example:



$M_2 = ?$
 $P_2 = ?$

$\theta + \omega = \omega_1$ everywhere (from m^+)

$\theta - \omega = \text{const} = 2\theta - \omega_1$ on m^-

$\Rightarrow \theta = \text{const. on } m^- \Rightarrow \omega = \text{const on } m^-$

\therefore characteristics
 are straight lines

$$M_1 = 1.775 \Rightarrow \omega_1 = 20 \text{ (from table)}$$

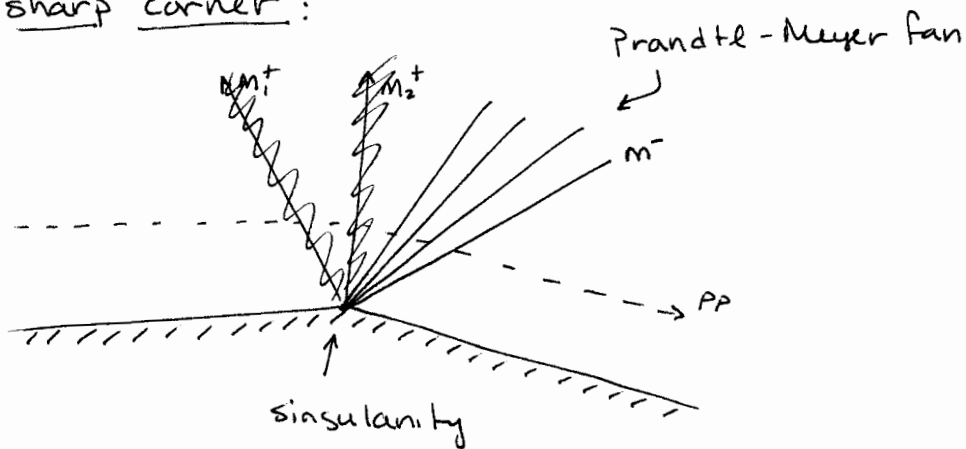
$$\omega = \omega_1 - \theta \text{ on } m^-$$

$$= 20 - (-16) = 36 \Rightarrow \boxed{M_2 = 2.369} \text{ (from table)}$$

$$\frac{P_1}{P_0} = \frac{0.18}{1.18} \text{ (from table)} \quad \textcircled{a} \quad M = 1.775 \Rightarrow P_0 = \frac{1}{1.18} \text{ atm}$$

$$\textcircled{a} \quad M_2 = 2.369 \quad \frac{P_2}{P_0} = \frac{0.0717}{0.18} = \boxed{0.397 \text{ atm}}$$

For a sharp corner:



on m^- $\theta = 0$, $M = 1.775$, $\omega = 20^\circ$; Sol'n same as above.

on m_2^+ $\theta = 16^\circ$, $M = 2.369$, $\omega = 36^\circ$