

6.730 Physics for Solid State Applications

Lecture 21:

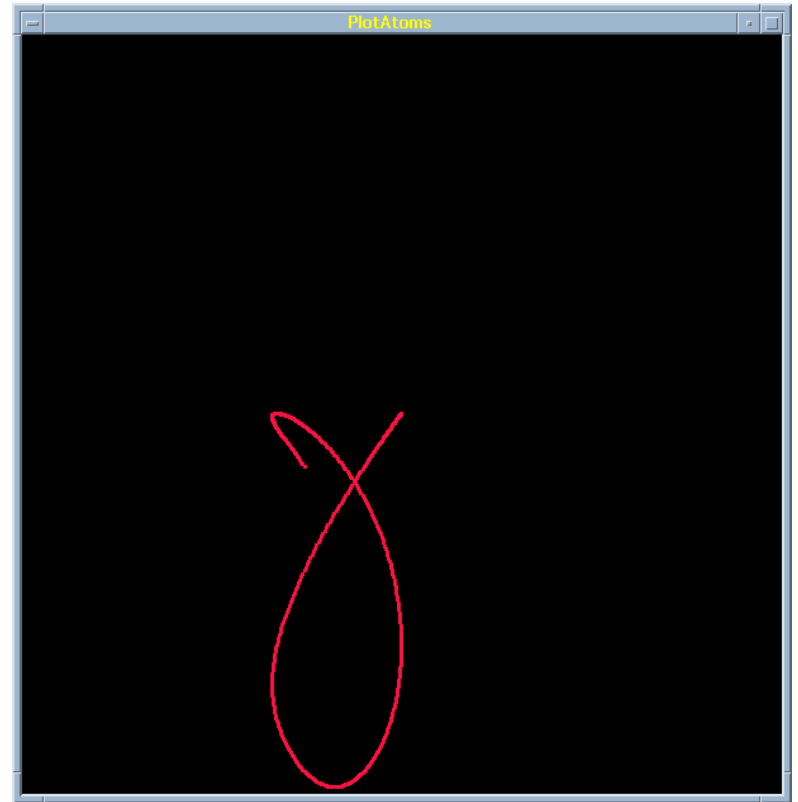
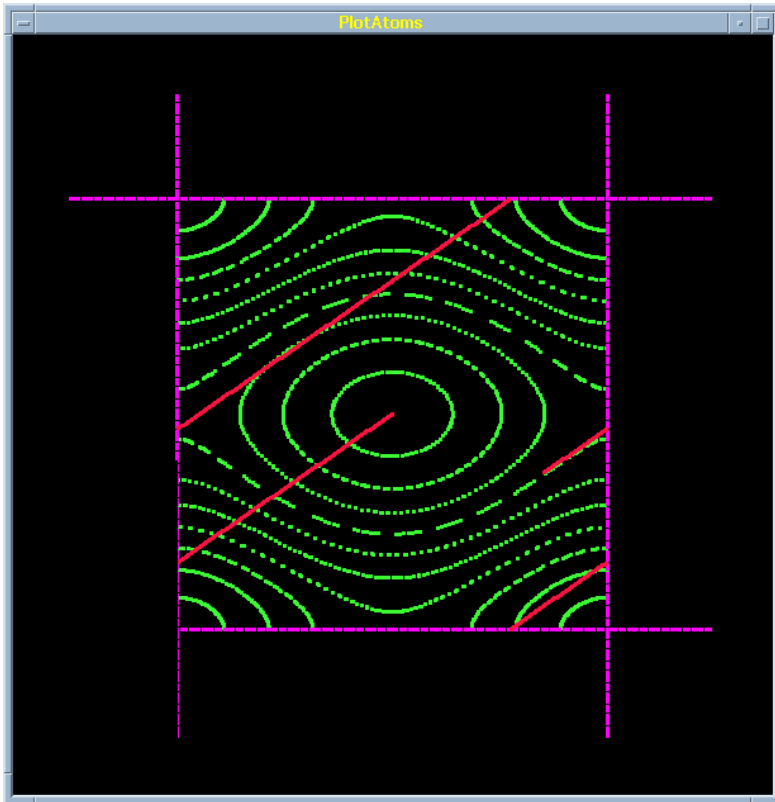
Outline

- Dynamical Effective Mass
- Fermi Surfaces
- Electrons and Holes

Semiclassical Equations of Motion

$$\langle \mathbf{v}_n(\mathbf{k}) \rangle = \frac{\langle \mathbf{p} \rangle}{m} = \frac{1}{\hbar} \nabla_{\mathbf{k}} E_n(\mathbf{k})$$

$$\mathbf{F}_{\text{ext}} = \hbar \frac{d\mathbf{k}}{dt}$$



Semiclassical Equations of Motion

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$$\mathbf{F}_{\text{ext}} = \hbar \frac{d\mathbf{k}}{dt}$$

Lets try to put these equations together....

$$\begin{aligned} a(t) &= \frac{dv}{dt} = \frac{1}{\hbar} \frac{\partial}{\partial t} \frac{\partial E_N(k)}{\partial k} = \frac{1}{\hbar} \frac{\partial^2 E_N(k)}{\partial k^2} \frac{dk}{dt} \\ &= \left[\frac{1}{\hbar^2} \frac{\partial^2 E_N(k)}{\partial k^2} \right] F_{\text{ext}} \end{aligned}$$

Looks like Newton's Law if we define the mass as follows...

$$m^*(k) = \hbar^2 \left(\frac{\partial^2 E_N(k)}{\partial k^2} \right)^{-1} \quad \text{dynamical effective mass}$$

 mass changes with k...so it changes with time according to k

Dynamical Effective Mass (3D)

Extension to 3-D requires some care,

\mathbf{F} and \mathbf{a} don't necessarily point in the same direction

$$\mathbf{a} = \overline{\overline{\mathbf{M}}}^{-1} \mathbf{F}_{\text{ext}} \quad \text{where} \quad \overline{\overline{\mathbf{M}}}_{i;j}^{-1} = \frac{1}{\hbar^2} \frac{\partial^2 E_N}{\partial k_i \partial k_j}$$

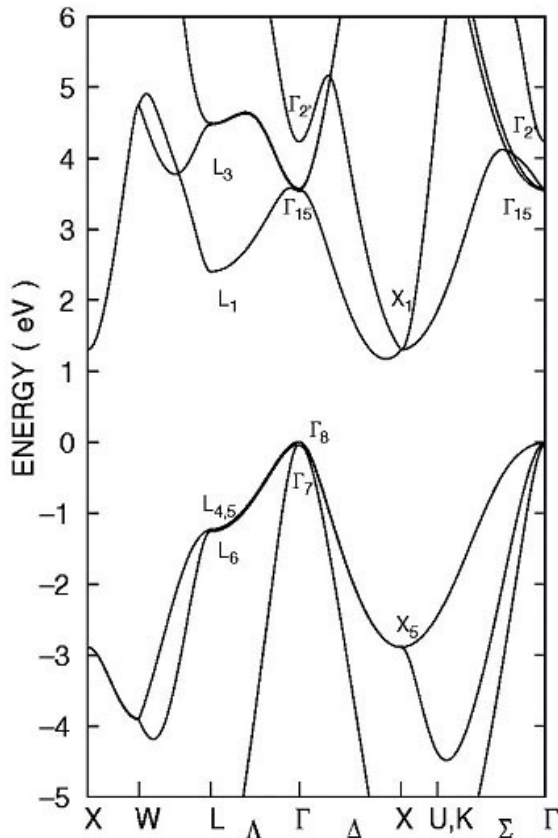
$$\begin{pmatrix} a_x \\ a_y \\ a_z \end{pmatrix} = \frac{d}{dt} \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} = \begin{pmatrix} \frac{1}{m_{xx}} & \frac{1}{m_{xy}} & \frac{1}{m_{xz}} \\ \frac{1}{m_{yx}} & \frac{1}{m_{yy}} & \frac{1}{m_{yz}} \\ \frac{1}{m_{zx}} & \frac{1}{m_{zy}} & \frac{1}{m_{zz}} \end{pmatrix} \begin{pmatrix} F_x \\ F_y \\ F_z \end{pmatrix}$$

Dynamical Effective Mass (3D)

Ellipsoidal Energy Surfaces

Fortunately, energy surfaces can often be approximate as...

$$E_N(k) = E_c + \frac{\hbar^2}{2} \left(\frac{(k_x - k_x^0)^2}{m_t} + \frac{(k_y - k_y^0)^2}{m_t} + \frac{(k_z - k_z^0)^2}{m_l} \right)$$



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WAVEVECTOR k

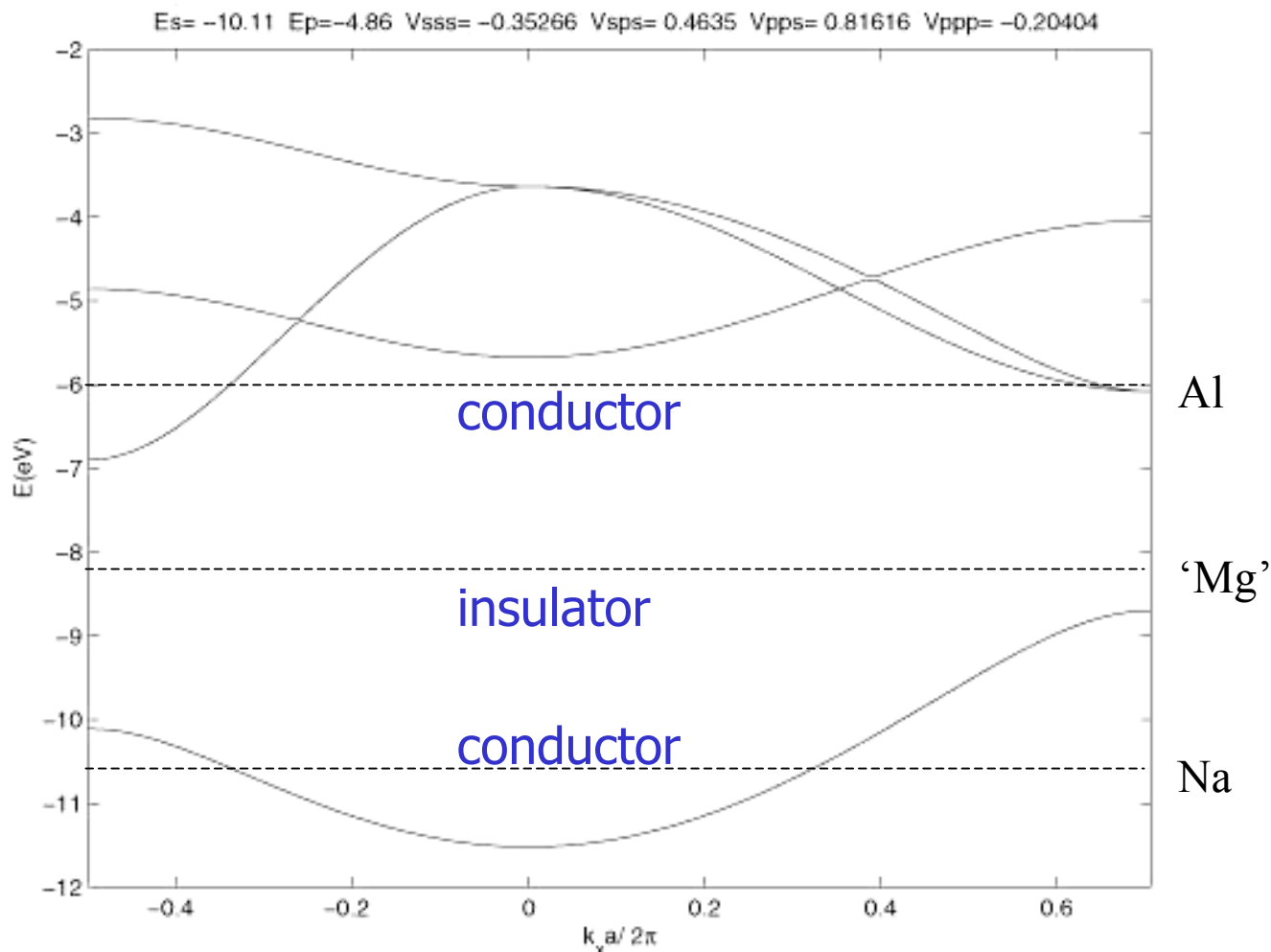
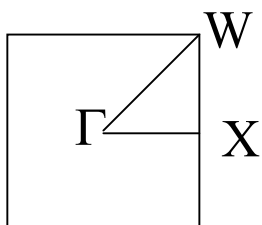
$$\overline{\overline{\mathbf{M}}}^{-1} = \begin{pmatrix} \frac{1}{m_t} & 0 & 0 \\ 0 & \frac{1}{m_t} & 0 \\ 0 & 0 & \frac{1}{m_l} \end{pmatrix}$$

$$\overline{\overline{\mathbf{M}}} = \begin{pmatrix} m_t & 0 & 0 \\ 0 & m_t & 0 \\ 0 & 0 & m_l \end{pmatrix}$$

2D Monatomic Square Crystals

Dispersion Relations

$a = 5.5 \text{ \AA}$



I	II	III	IV	V	VI	VII	VIII
11	12	13	14	15	16	17	18
Na	Mg	Al	Si	P	S	Cl	Ar

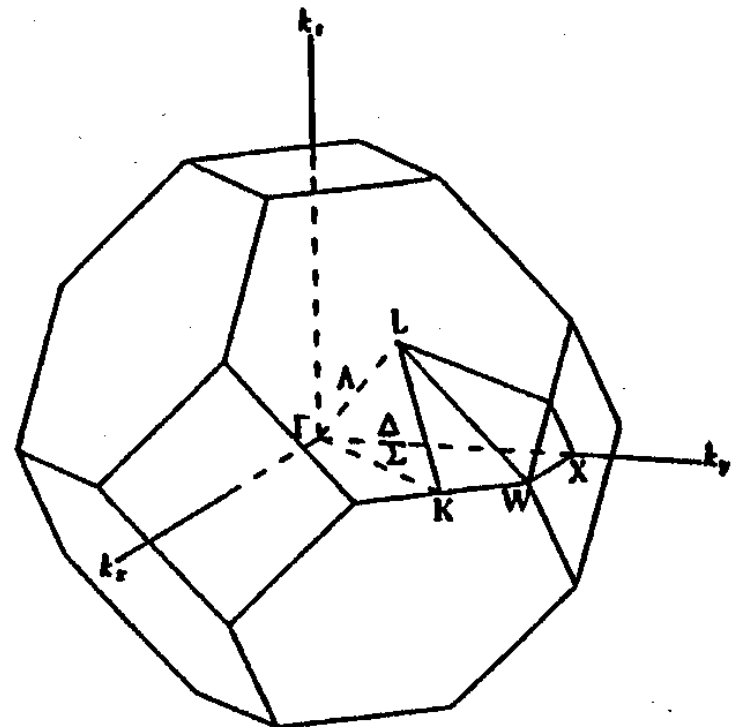
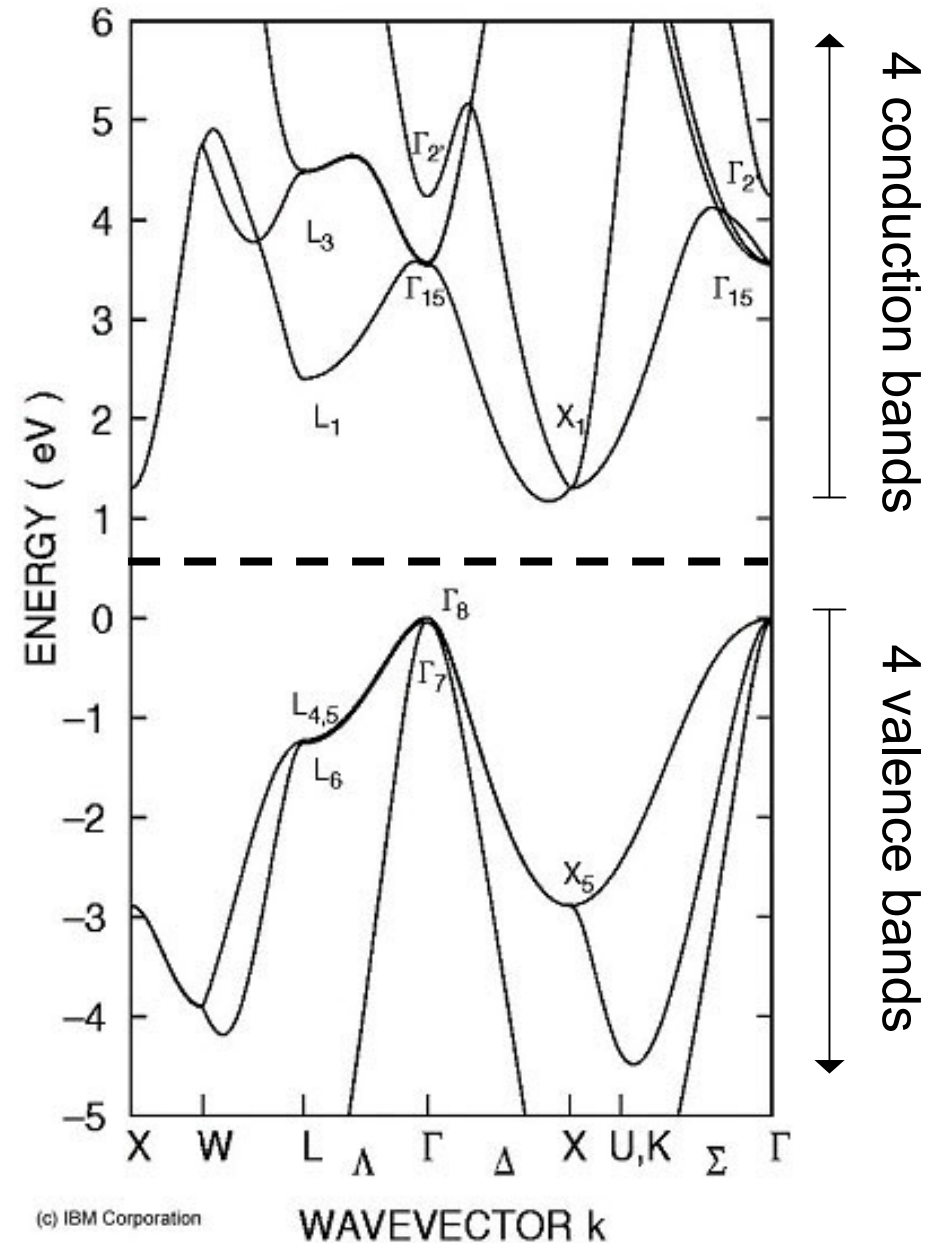
Silicon Bandstructure



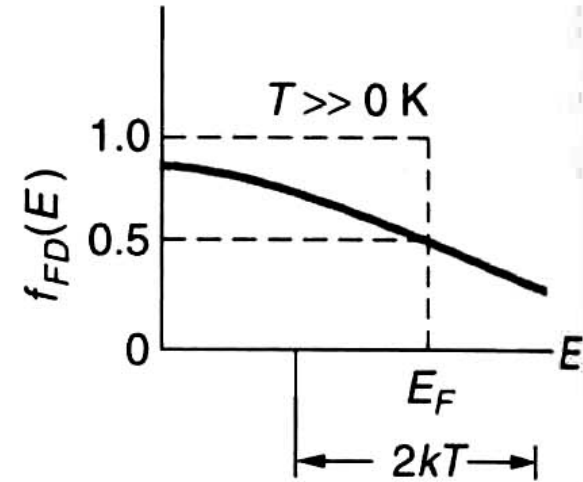
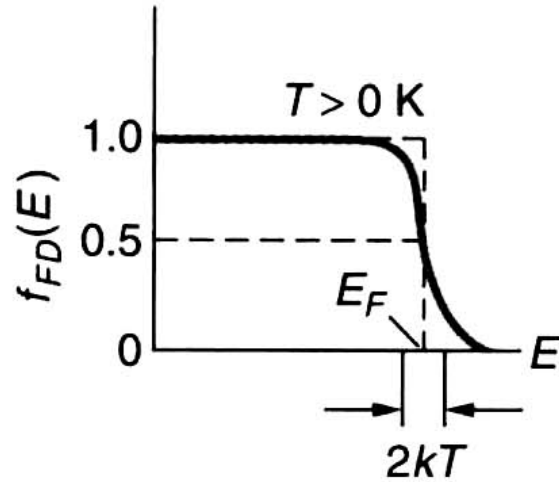
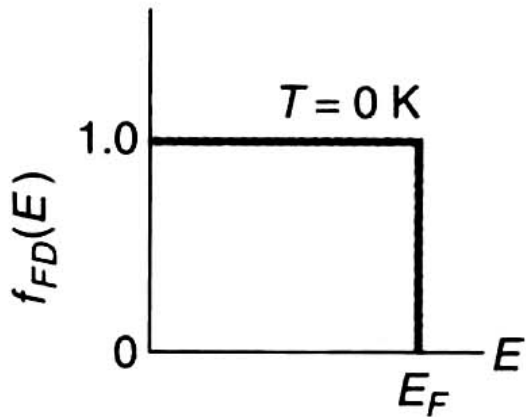
4 e- per silicon atom

2 silicon atoms per lattice site

total: 8 electrons at each site



Finite Temperatures



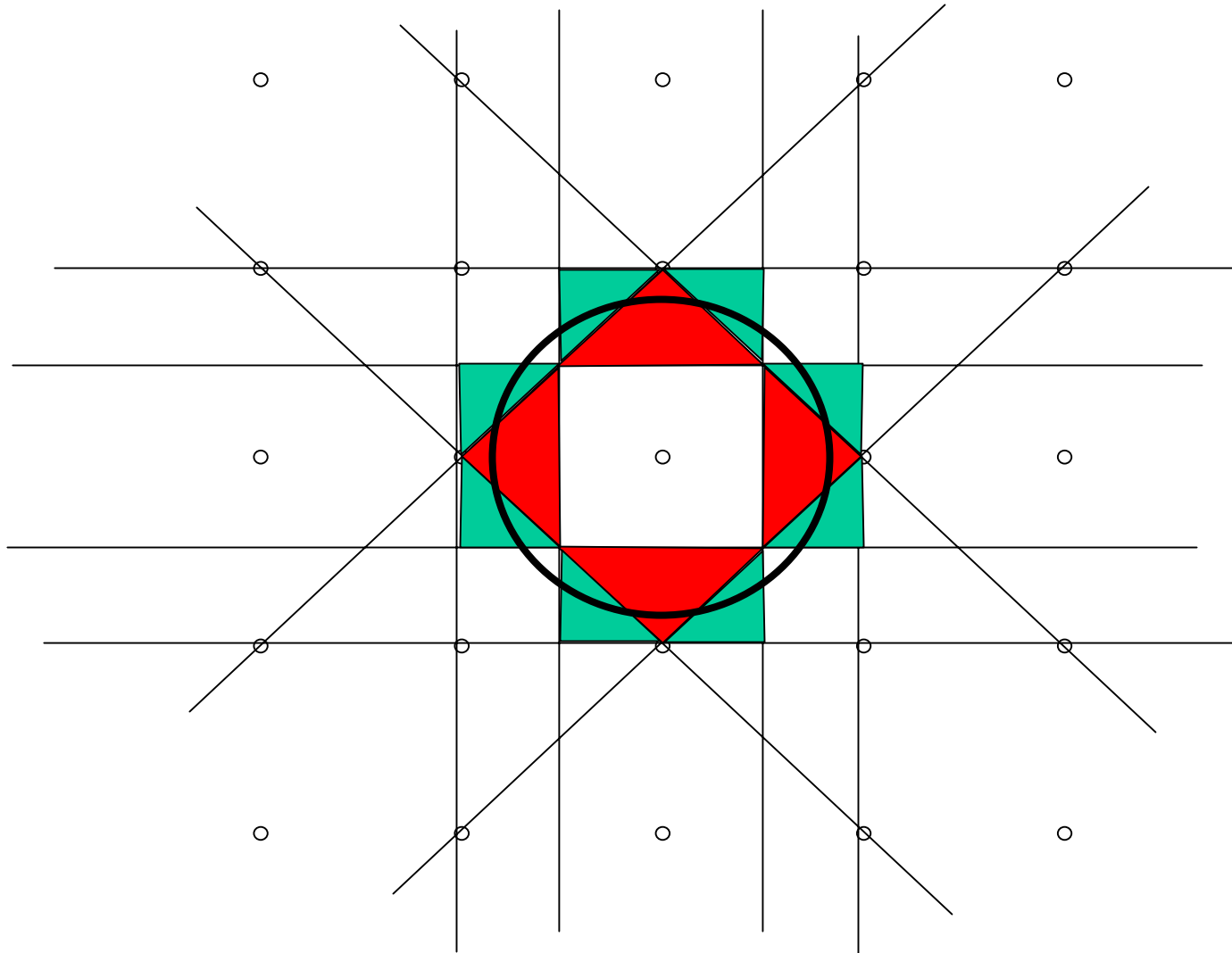
$$n = \frac{N}{V} = \int_{-\infty}^{\infty} \frac{1}{1 + e^{(E_{\mathbf{k}} - \mu)/k_B T}} 2 \frac{d^3 \mathbf{k}}{(2\pi)^3}$$

$$\frac{E_{\text{total}}}{V} = \int_{-\infty}^{\infty} E_{\mathbf{k}} \frac{1}{1 + e^{(E_{\mathbf{k}} - \mu)/k_B T}} 2 \frac{d^3 \mathbf{k}}{(2\pi)^3}$$

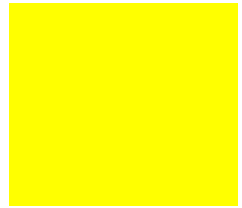
Free Electron Fermi Surfaces (2D)

T=0

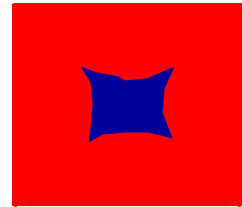
For free electrons energy surfaces are simple spheres (circles)...
Valence (# of electrons) determines radius of energy surface...



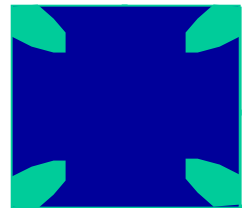
1st zone



2nd zone



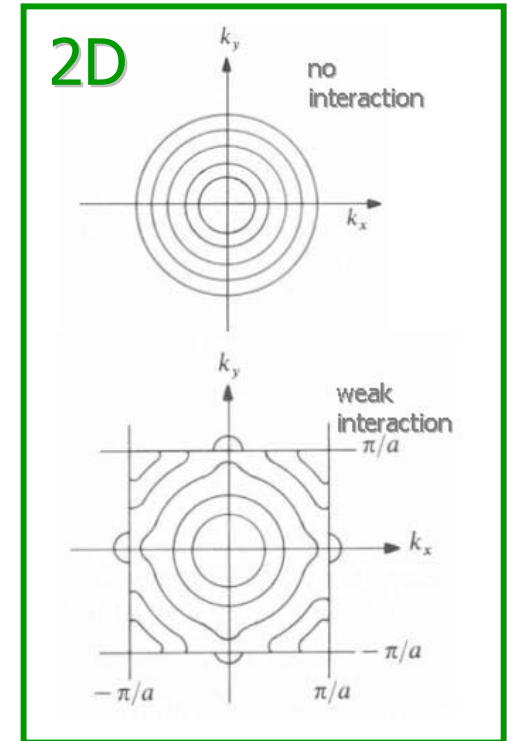
3rd zone



Fermi Surfaces (3D)

When k near to BZ boundary:

→ E contours become distorted

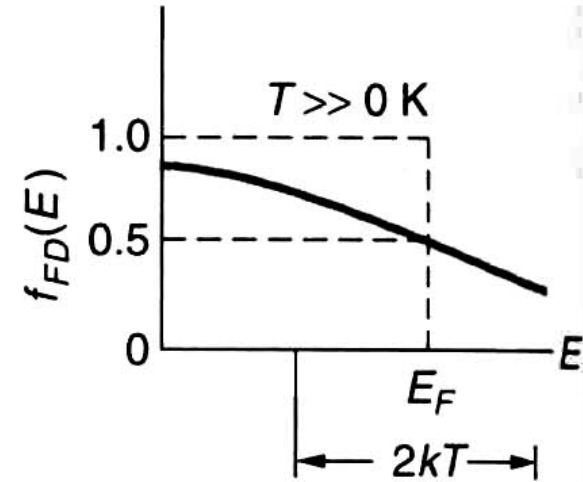
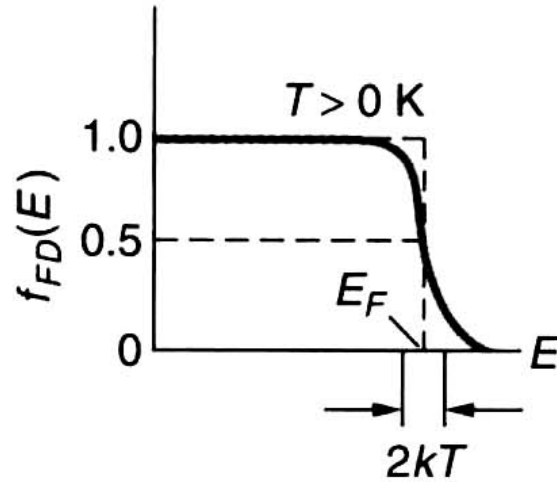
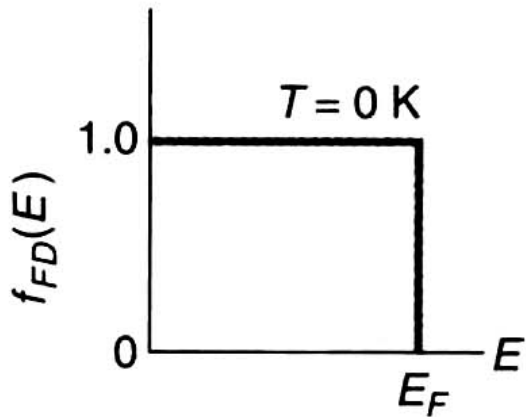


Fermi Surfaces (3D)

$N_e = 1$ monovalent metals, e.g. Na, Cu, with values \sim f.e. theory

other cases, e.g. Be ($N_e=2$), Al ($N_e=3$), there are serious differences

Finite Temperatures



$$n = \frac{N}{V} = \int_{-\infty}^{\infty} \frac{1}{1 + e^{(E_{\mathbf{k}} - \mu)/k_B T}} 2 \frac{d^3 \mathbf{k}}{(2\pi)^3}$$

$$\frac{E_{\text{total}}}{V} = \int_{-\infty}^{\infty} E_{\mathbf{k}} \frac{1}{1 + e^{(E_{\mathbf{k}} - \mu)/k_B T}} 2 \frac{d^3 \mathbf{k}}{(2\pi)^3}$$

Overview of Electron Distributions



Metal

Insulator

or

Semiconductor

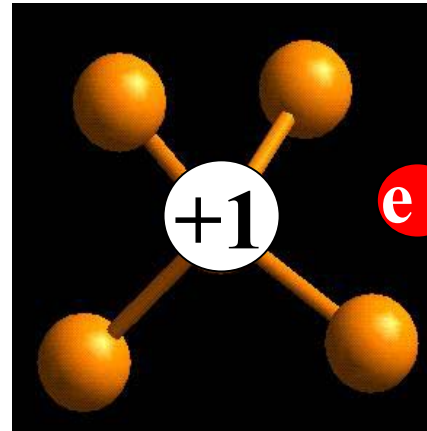
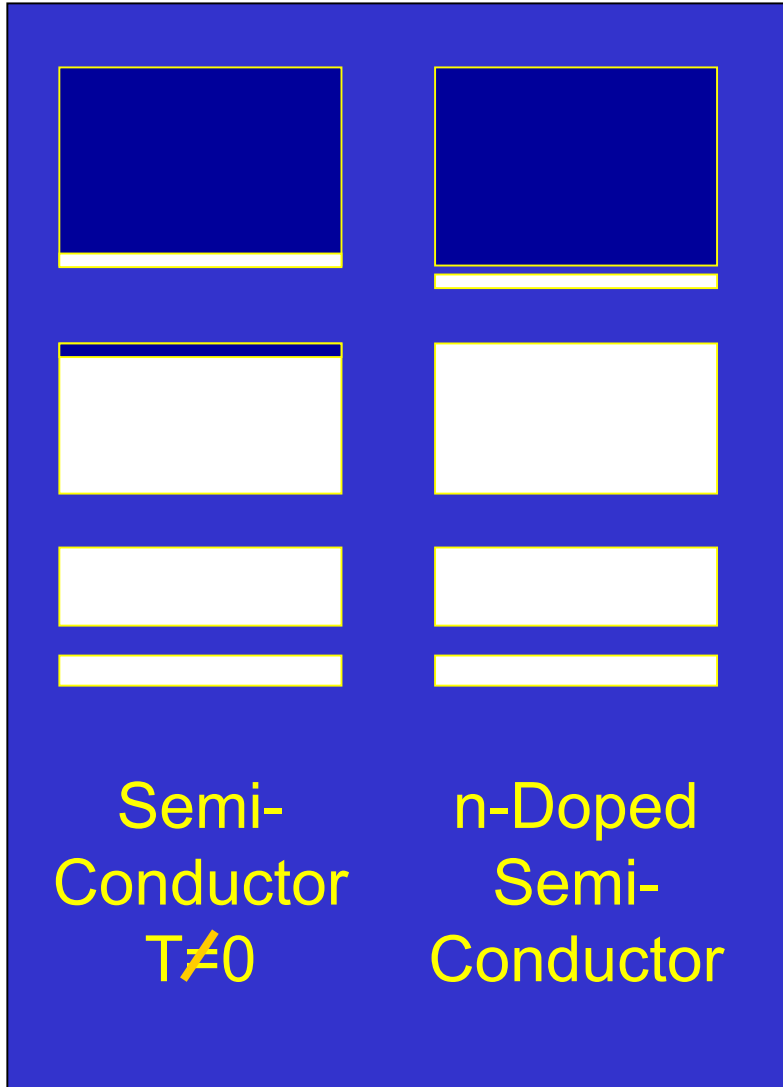
$T=0$

Semi-
Conductor

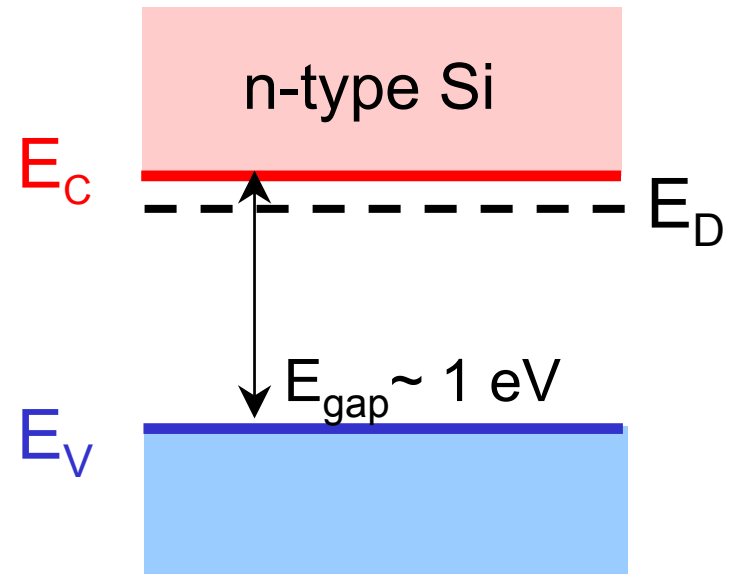
$T \neq 0$

n-Doped
Semi-
Conductor

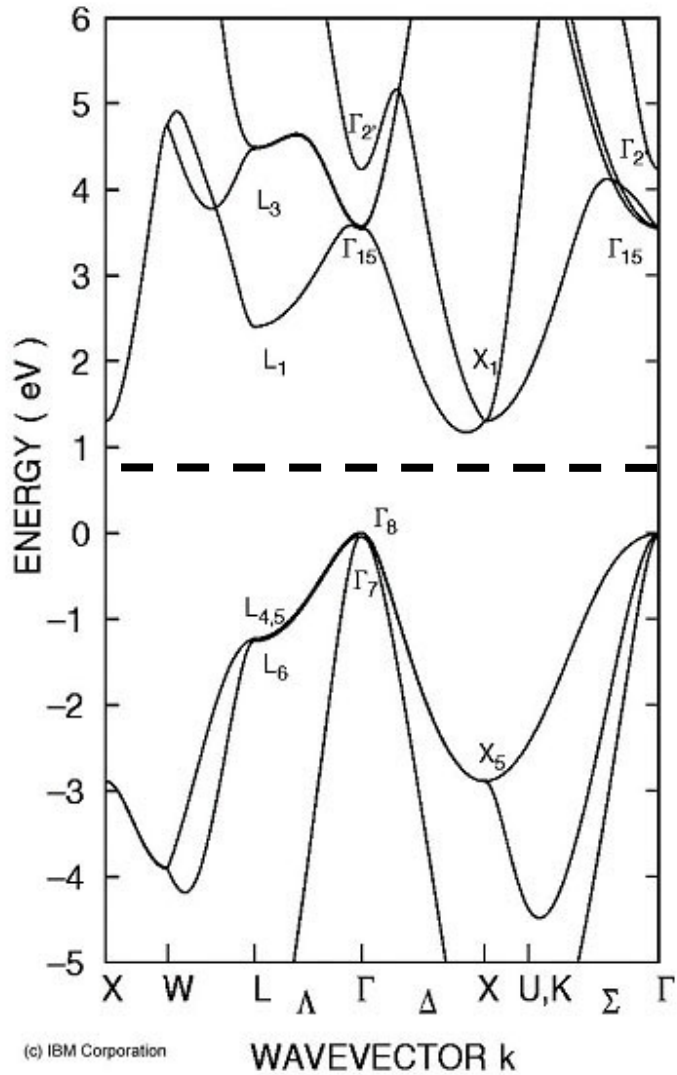
Electron Distributions in Doped Semiconductors



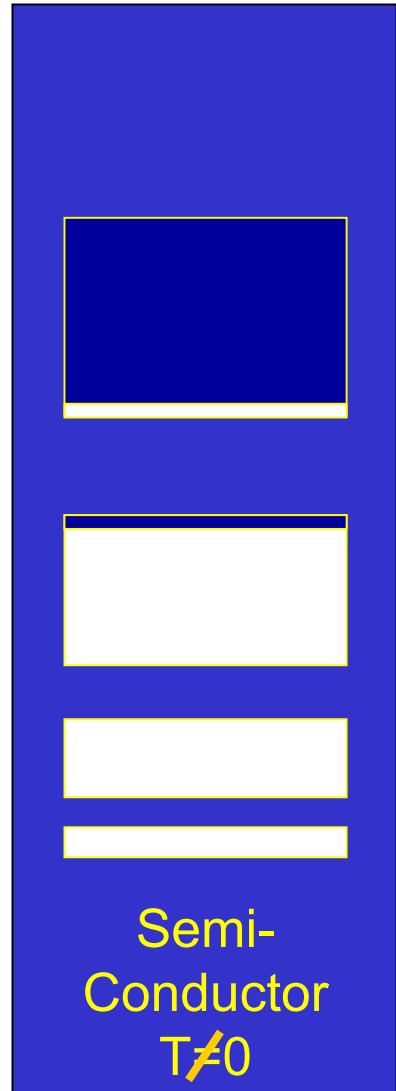
$$E_d = E_c - \frac{13.56 m^*}{l^2 \epsilon^2} \frac{m}{m} \text{ eV}$$



Electron and Holes



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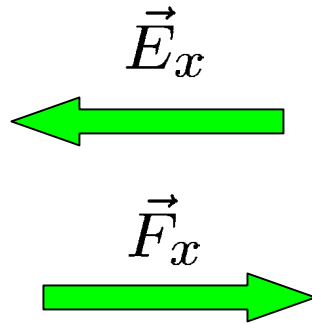


Electrons in conduction band

Holes in valence band

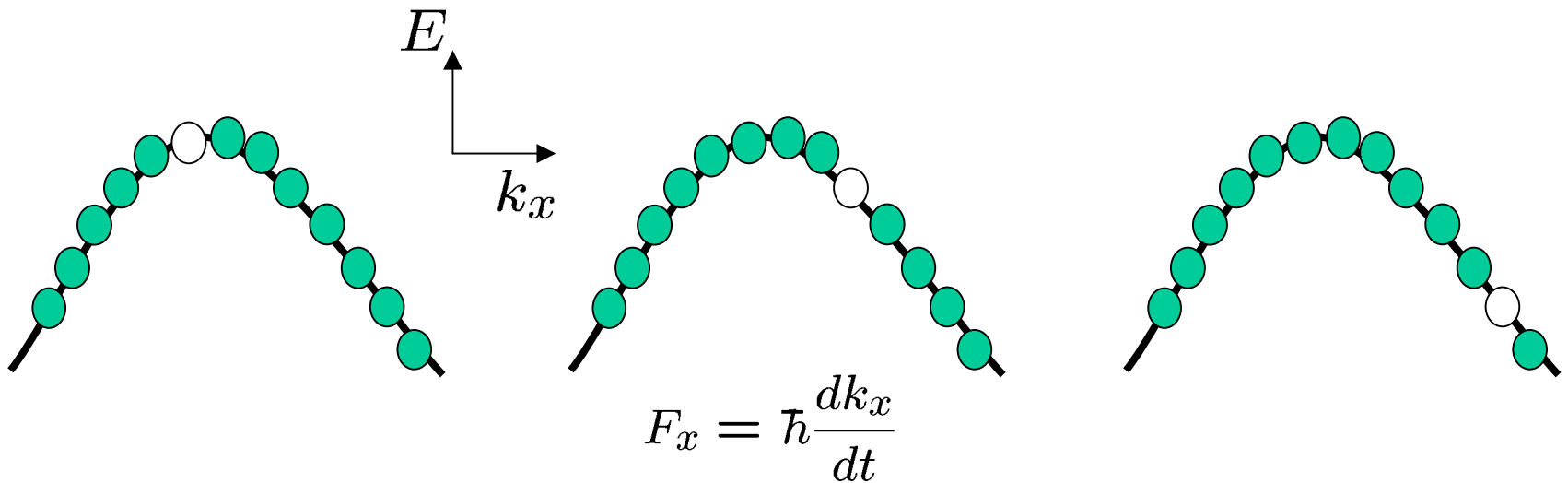
Semi-
Conductor
 $T \neq 0$

Motion of Valence Electrons



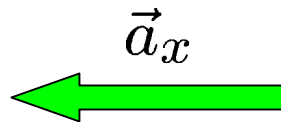
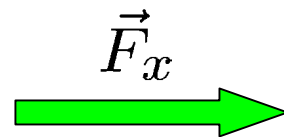
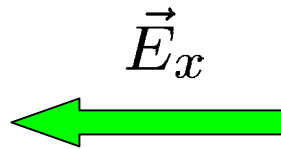
electrons have negative charge

k-space



Valence electrons (and vacancy) all move in the positive k_x direction...

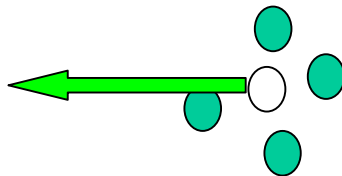
Motion of Valence Electrons



electrons have negative charge

valence electrons have negative mass !

Real space



Vacancy ends up moving in the direction of the electric field as if it had a positive charge

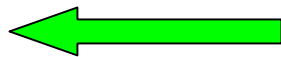
Hole is a quasi-particle with positive charge and positive mass...

Motion of Valence Electrons

$$\vec{E}_x$$

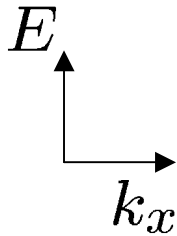


$$\vec{a}_x$$



Hole is a quasi-particle with positive charge and positive mass...

k-space



hole dispersion

