

MIT OpenCourseWare
<http://ocw.mit.edu>

6.453 Quantum Optical Communication
Spring 2009

For information about citing these materials or our Terms of Use, visit: <http://ocw.mit.edu/terms>.



October 16, 2008

6.453 *Quantum Optical Communication* Lecture 11

Jeffrey H. Shapiro

Optical and Quantum Communications Group

RESEARCH LABORATORY OF ELECTRONICS
Massachusetts Institute of Technology

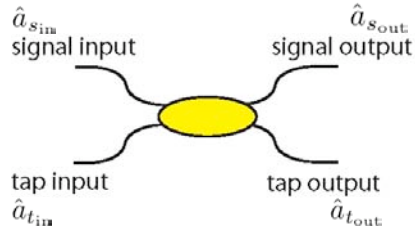
www.rle.mit.edu/qoptics

6.453 *Quantum Optical Communication* - Lecture 11

- Announcements
 - Pick up lecture notes, slides
- Single-Mode Photodetection
 - Squeezed-state waveguide tap — reprise
- Single-Mode Linear Systems
 - Attenuators
 - Phase-Insensitive Amplifiers

Optical Waveguide Tap — Quantum

Fused Fiber Coupler



- Coupler is a beam splitter

$$\hat{a}_{s_{out}} = \sqrt{T}\hat{a}_{s_{in}} + \sqrt{1-T}\hat{a}_{t_{in}}$$

$$\hat{a}_{t_{out}} = \sqrt{1-T}\hat{a}_{s_{in}} - \sqrt{T}\hat{a}_{t_{in}}$$

- Tap input is in squeezed vacuum
- Homodyne SNR at signal input

$$\text{SNR}_{in} = 4|a_{s_{in}}|^2$$

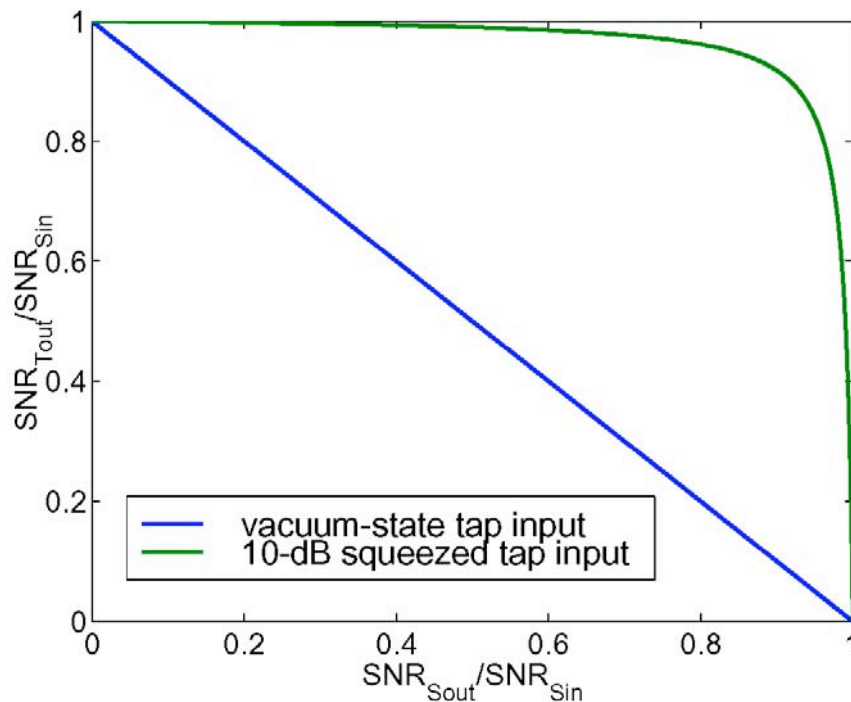
- Homodyne SNR at signal output

$$\text{SNR}_{out} = \frac{4T|a_{s_{in}}|^2}{T + (1-T)(\mu - \nu)^2}$$

- Homodyne SNR at tap output

$$\text{SNR}_{tap} = \frac{4(1-T)|a_{s_{in}}|^2}{(1-T) + T(\mu - \nu)^2}$$

Optical Waveguide Tap: SNR Tradeoff



Non-Ideal Quantum Photodetection

- Quantum Efficiency $\eta < 1$:

$$\hat{a}' \equiv \sqrt{\eta} \hat{a} + \sqrt{1 - \eta} \hat{a}_\eta, \quad \hat{a}_\eta \text{ in vacuum state}$$

- Direct Detection: $\hat{a}'^\dagger \hat{a}'$ measurement
- Homodyne Detection: $\text{Re}(\hat{a}' e^{-j\theta})$ measurement
- Heterodyne Detection: \hat{a}' measurement

Non-Ideal Quantum Photodetection: $\eta < 1$

- Direct Detection of Number State $|n\rangle$:

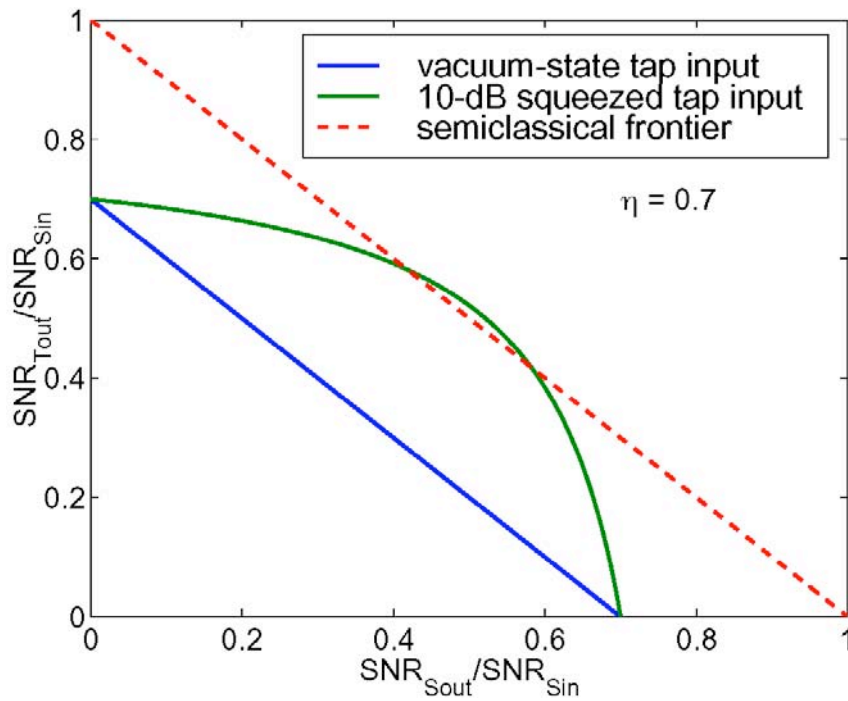
$$\text{Pr}(\hat{N}' = k \mid |n\rangle) = \binom{n}{k} \eta^k (1 - \eta)^{n-k}, \quad \text{for } 0 \leq k \leq n$$

- Homodyne Detection of Squeezed State $|\beta; \mu, \nu\rangle$:

$$\alpha'_1 \sim N\left(\sqrt{\eta}(\mu - \nu)\beta, \frac{\eta(\mu - \nu)^2 + (1 - \eta)}{4}\right),$$

for β, μ, ν real

Loss is a Problem for Squeezed States



Single-Mode Linear Systems: the Attenuator

- Classical Attenuation

$$a_{\text{in}} \longrightarrow \boxed{0 < L < 1} \longrightarrow a_{\text{out}} = \sqrt{L} a_{\text{in}}$$

- Noiseless attenuation is possible

$$\frac{\langle |a_{\text{out}}|^2 \rangle^2}{\langle \Delta(|a_{\text{out}}|^2) \rangle^2} = \frac{\langle |a_{\text{in}}|^2 \rangle^2}{\langle \Delta(|a_{\text{in}}|^2) \rangle^2}$$

$$\frac{\langle a_{\text{out}\theta} \rangle^2}{\langle \Delta a_{\text{out}\theta}^2 \rangle} = \frac{\langle a_{\text{in}\theta} \rangle^2}{\langle \Delta a_{\text{in}\theta}^2 \rangle}$$

Single-Mode Linear Systems: the Amplifier

- Classical Amplification

$$a_{\text{in}} \longrightarrow \boxed{G > 1} \longrightarrow a_{\text{out}} = \sqrt{G} a_{\text{in}}$$

- Noiseless amplification is possible

$$\frac{\langle |a_{\text{out}}|^2 \rangle^2}{\langle \Delta(|a_{\text{out}}|^2) \rangle^2} = \frac{\langle |a_{\text{in}}|^2 \rangle^2}{\langle \Delta(|a_{\text{in}}|^2) \rangle^2}$$

$$\frac{\langle a_{\text{out}\theta} \rangle^2}{\langle \Delta a_{\text{out}\theta}^2 \rangle} = \frac{\langle a_{\text{in}\theta} \rangle^2}{\langle \Delta a_{\text{in}\theta}^2 \rangle}$$

Single-Mode Linear Systems: Quantum Case

- Input and Output are Photon Annihilation Operators:

$$\left[\hat{a}_{\text{out}}, \hat{a}_{\text{out}}^\dagger \right] = \left[\hat{a}_{\text{in}}, \hat{a}_{\text{in}}^\dagger \right] = 1$$

- Quantum Attenuator *Cannot* Obey $\hat{a}_{\text{out}} = \sqrt{L} \hat{a}_{\text{in}}$
- Quantum Amplifier *Cannot* Obey $\hat{a}_{\text{out}} = \sqrt{G} \hat{a}_{\text{in}}$
- Loss and Gain Require Additional Quantum Noise:
to preserve the Heisenberg Uncertainty Principle

Single-Mode Linear Systems: Quantum Case

- Quantum Attenuation

$$\hat{a}_{\text{in}} \longrightarrow \boxed{0 < L < 1} \longrightarrow \hat{a}_{\text{out}} = \sqrt{L} \hat{a}_{\text{in}} + \sqrt{1-L} \hat{a}_L$$

- Quantum Amplification

$$\hat{a}_{\text{in}} \longrightarrow \boxed{G > 1} \longrightarrow \hat{a}_{\text{out}} = \sqrt{G} \hat{a}_{\text{in}} + \sqrt{G-1} \hat{a}_G^\dagger$$

Single-Mode Linear Systems: Quantum Case

- Attenuator with Coherent-State \hat{a}_{in} and Vacuum-State \hat{a}_L :

$$\langle \Delta \hat{N}_{\text{out}}^2 \rangle = L \langle \hat{N}_{\text{in}} \rangle$$

$$\langle \Delta \hat{a}_{\text{out}\theta}^2 \rangle = \langle \Delta \hat{a}_{\text{in}\theta}^2 \rangle = \frac{1}{4}$$

- Amplifier with Coherent-State \hat{a}_{in} and Vacuum-State \hat{a}_G :

$$\langle \Delta \hat{N}_{\text{out}}^2 \rangle =$$

$$[G \langle \hat{N}_{\text{in}} \rangle + (G-1)] + [2G(G-1) \langle \hat{N}_{\text{in}} \rangle + (G-1)^2]$$

$$\langle \Delta \hat{a}_{\text{out}\theta}^2 \rangle = \frac{2G-1}{4}$$

Coming Attractions: Lecture 12

- Lecture 12:
Single-Mode and Two-Mode Linear Systems
 - Phase-Insensitive Amplifiers
 - Phase-Sensitive Amplifiers
 - Entanglement