

TIME-CORRELATION FUNCTIONS

- > T.C.F.s are statistical measures of the time-evolution of an observable under equilibrium conditions. → They describe an ensemble.
- > T.C.F.s can be used to describe spectroscopy and other time-dependent phenomena.
- > They are generally applicable to any time-dependent process for an ensemble, but are commonly used to describe random or stochastic processes in condensed phases.

Qualitatively:

- A correlation function describes how long a given property of a system persists until it is averaged out by microscopic motions of system.
- It describes how and when a statistical relationship has vanished.
- T.C.F.s describes the dynamics associated with a dynamical variable $A(t)$, averaged over the ensemble.

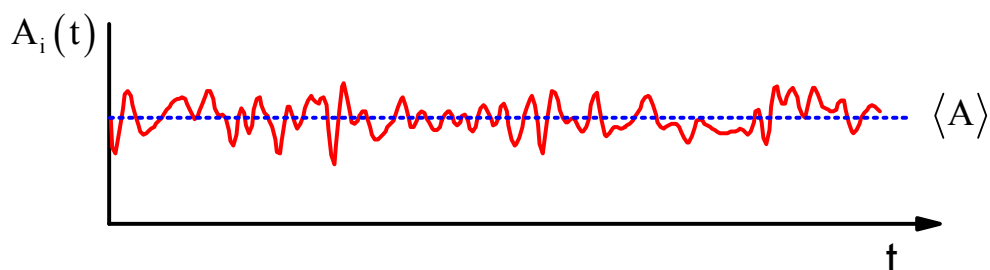
A : Microscopic variable

$\langle A \rangle$: Equilibrium ensemble average

Classical: $\langle A \rangle \equiv \int d\bar{p} \int d\bar{q} A(p, q; t) f(p, q)$ f : equil. Prob. distribution function

Quantum: $\langle A \rangle = \sum_n p_n \langle n | A | n \rangle$ $p_n = e^{-\beta E_n} / Z$

$A_i(t)$: Expectation value of A for a member of ensemble as a function of time.



If we look at this behavior there seems to be little information in the random fluctuations of A , but there are characteristic time scales and amplitudes to these changes. We can characterize these by comparing the value of A at time t with the value of A at time t' later.

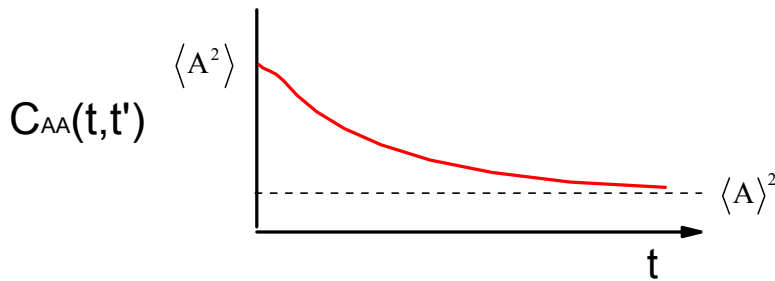
Correlation functions are defined as a time-dependent quantity, $A(t)$, multiplied by that quantity at some later time, $A(t')$, and averaged over ensemble:

$$C_{AA}(t, t') \equiv \langle A(t)A(t') \rangle \quad \leftarrow \text{auto-correlation function}$$

$$C_{AB}(t, t') \equiv \langle A(t)B(t') \rangle \quad \leftarrow \text{cross-correlation function}$$

These are products of a pair of dynamical variables.

Properties of Correlation Functions



$$1. \quad C_{AA}(t, t) = \langle A(t)A(t) \rangle = \langle A^2(t) \rangle \geq 0 \\ = \langle A^2 \rangle \quad (\text{mean square value of } A - \text{independent of time})$$

$$2. \quad C_{AA}(t, t') = C_{AA}(t + \tau, t' + \tau) = C_{AA}(t - t') \quad \leftarrow \text{for } \tau = -t'$$

↗ Stationary random process.
Time average = Ensemble average

↖ Just express in time separation $C(t)$

$$3. \quad \lim_{t \rightarrow \infty} C_{AA}(t) = \langle A(t) \rangle \langle A(0) \rangle = \langle A \rangle^2$$

Processes lose all correlation at infinite time separation .

4. For classical mechanics:

$$\langle A(t)A(t') \rangle = \langle A(t')A(t) \rangle$$

$$C_{AA}(t) = C_{AA}(-t) \quad \text{Even in time}$$

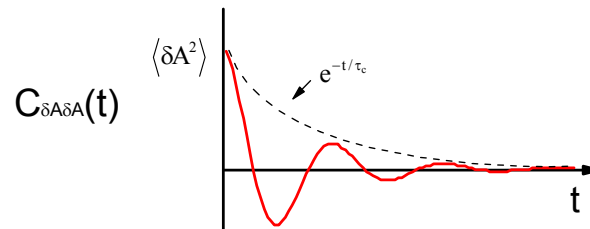
5. If we are observing fluctuations about an average position:

Redefine $\delta A \equiv A - \langle A \rangle$ deviation from average

$$C_{\delta A \delta A}(t) = \langle \delta A(t) \delta A(0) \rangle = C_{AA}(t) - \langle A \rangle^2$$

Properties: $\lim_{t \rightarrow \infty} C_{\delta A \delta A}(t) = 0$ all correlation lost at $t = \infty$

$$C_{\delta A \delta A}(0) = \langle \delta A^2 \rangle = \langle A^2 \rangle - \langle A \rangle^2 \quad (\text{variance})$$



6. τ_c : Correlation time is characteristic time for decay to zero.

There may be several time scales over which the correlation vanishes, so we can define

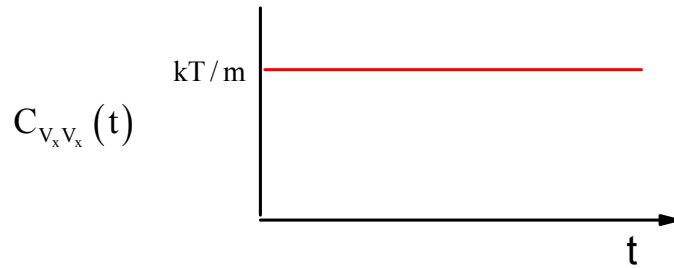
$$\tau_c = \frac{1}{\langle \delta A^2 \rangle} \int_0^{\infty} dt \langle \delta A(t) \delta A(0) \rangle$$

EXAMPLE 1: Velocity autocorrelation function for gas.

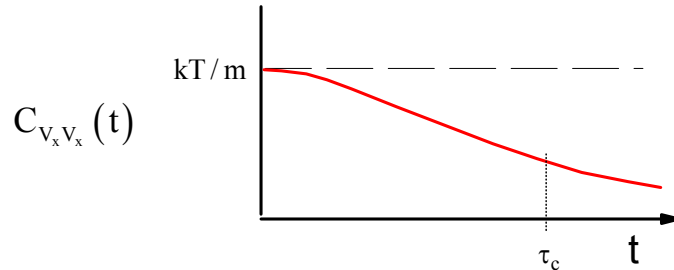
V_x : \hat{x} Component of molecular velocity $\langle V_x \rangle = 0$

$C_{V_x V_x}(t) = \langle V_x(t) V_x(0) \rangle$ $C_{V_x V_x}(0) = \langle V_x^2(0) \rangle = \frac{kT}{m}$

Ideal gas: no collisions



Dilute gas: infrequent collisions



$V_x(t) = V_x(0)$ for $t < \tau_c$
 $V_x(t) = V_x(0) \pm \delta$ for $t \geq \tau_c$

τ_c is related to the mean time between collisions.

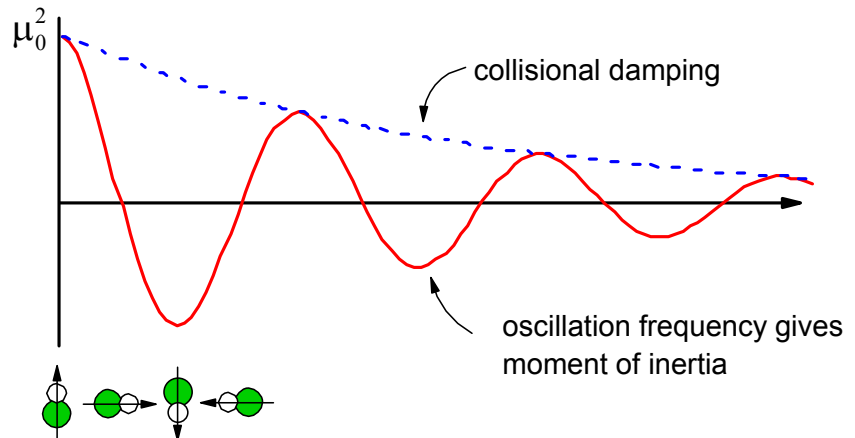
EXAMPLE 2: $\bar{\mu}_i$: dipole moment for diatomic molecule in dilute gas.

$\langle \bar{\mu}_i \rangle = 0$ (all angles are equally likely: isotropic system)

$\bar{\mu}_i = \mu_0 \cdot \hat{u}$
↖ unit vector along dipole

$C_{\mu\mu}(t) = \langle \bar{\mu}(t) \bar{\mu}(0) \rangle$
 $= \mu_0^2 \langle \hat{u}(t) \cdot \hat{u}(0) \rangle$

Project time-dependent orientation onto initial orientation



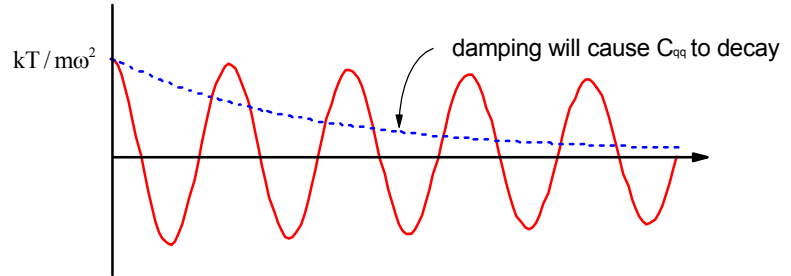
EXAMPLE 3: Displacement of harmonic oscillator.

$$m\ddot{q} = -\kappa q \quad \rightarrow \quad \ddot{q} = -\omega^2 q$$

$$q(t) = q(0) \cos \omega t$$

$$\text{Since } \langle q^2(0) \rangle = \frac{kT}{m\omega^2}$$

$$\begin{aligned} C_{qq}(t) &= \langle q(t)q(0) \rangle = \langle q^2(0) \rangle \cos \omega t \\ &= \left(\frac{kT}{m\omega^2} \right) \cos \omega t \end{aligned}$$



QUANTUM CORRELATION FUNCTIONS

Equilibrium (thermal) ensemble average over product of Hermetian operators evaluated two times.

→ Heisenberg Representation

$$C_{AA}(t, t') = \langle\langle A(t)A(t') \rangle\rangle$$

Here I'm temporarily using a double bracket to remind you that this is the equilibrium ensemble average over the expectation value of the product of operators. It is also sometimes written

$$\langle A(t)A(0) \rangle_{\text{eq}}$$

Referring back to our notation for mixed states, we can represent the state of the system through

$$|\Psi_k\rangle = \sum_n a_n^k |n\rangle$$

$$\begin{aligned} C_{AA}(t, t') &= \sum_k \langle \Psi_k | A(t)A(t') | \Psi_k \rangle p_k \\ &= \sum_{k,n,m} p_k (a_n^k)^* a_m^k \langle n | A(t)A(t') | m \rangle \\ &= \sum_{n,m} \langle a_n^* a_m \rangle \langle n | A(t)A(t') | m \rangle \end{aligned}$$

The equilibrium ensemble average of the expansion coefficients is equivalent to phase averaging over the expansion coefficients, since at equilibrium all phases are equally probable:

$$\langle a_n^* a_m \rangle = \frac{1}{2\pi} \int_0^{2\pi} a_n^* a_m d\phi = \frac{1}{2\pi} |a_n| |a_m| \int_0^{2\pi} e^{-i(\phi_n - \phi_m)} d\phi_{nm} \quad \text{where } a_n = |a_n| e^{i\phi_n}$$

The integral is quite clearly zero unless $\phi_n = \phi_m$, giving

$$\langle a_n^* a_m \rangle = p_n \delta_{n,m} = \frac{e^{-\beta E_m}}{Z} \delta_{m,n}$$

So

$$C_{AA}(t, t') = \sum_n p_n \langle n | A(t) A(t') | n \rangle \quad \text{Heisenberg representation}$$

More commonly, we will be writing correlation functions in the interaction picture operators.

We can also write this in the Schrödinger picture:

$$U = \exp(-iHt/\hbar)$$

$$\begin{aligned} C_{AA}(t, t') &= \sum_n p_n \langle n | \underline{U^\dagger(t) A U(t) U^\dagger(t') A U(t') A} | n \rangle \\ &= \sum_n p_n \langle n | A U(t-t') A | n \rangle e^{+i\omega_n(t-t')} \\ &= \sum_{n,j} p_n \langle n | A | j \rangle \langle j | A | n \rangle e^{-i\omega_j(t-t')} \quad \text{insert } \sum_j |j\rangle \langle j| \\ &= \sum_{n,j} p_n |A_{jn}|^2 e^{-i\omega_j(t-t')} \end{aligned}$$

Notice that $C_{AA}(t, t') = C_{AA}(t-t') \Rightarrow C_{AA}(t)$

$$C_{AA}(-t) = C_{AA}^*(t)$$

Properties of Quantum Correlation Functions

There are various properties of quantum correlation functions that can be obtained using the properties of the time-evolution operator.

$$\begin{aligned}\langle A(0)A(t) \rangle &= \langle A(0)U^\dagger A U \rangle \\ &= \langle U A U^\dagger A \rangle \\ &= \langle A(-t)A(0) \rangle\end{aligned}$$

$$\begin{aligned}\langle A(t)A(0) \rangle^* &= \langle U^\dagger A U A \rangle^* \\ &= \langle U A U^\dagger A \rangle \\ &= \langle A(0)A(t) \rangle\end{aligned}$$

$$\therefore \langle A(-t)A(0) \rangle = \langle A(t)A(0) \rangle^* = \langle A(0)A(t) \rangle \qquad C_{AA}^*(t) = C_{AA}(-t)$$

$$\begin{aligned}\langle A(t)A(t') \rangle &= \langle U^\dagger(t)A(0)U(t)U^\dagger(t')A(0)U(t') \rangle \\ &= \langle U(t')U^\dagger(t)A U(t)U^\dagger(t')A \rangle \\ &= \langle U^\dagger(t-t')A U(t-t')A \rangle \\ &= \langle A(t-t')A(0) \rangle\end{aligned}$$

Note that $C_{AA}(t)$ is complex. You cannot directly measure a quantum correlation function, but observables are often related to the real or imaginary part of correlation functions, or other combinations of correlation functions.

$$C_{AA}(t) = C'_{AA}(t) + iC''_{AA}(t)$$

$$C'_{AA}(t) = \frac{1}{2}[C_{AA}(t) + C_{AA}^*(t)] = \frac{1}{2}[\langle A(t)A(0) \rangle + \langle A(0)A(t) \rangle] = \frac{1}{2}\langle [A(t), A(0)]_+ \rangle$$

$$C''_{AA}(t) = \frac{1}{2i}[C_{AA}(t) - C_{AA}^*(t)] = \frac{1}{2i}[\langle A(t)A(0) \rangle - \langle A(0)A(t) \rangle] = \frac{1}{2i}\langle [A(t), A(0)]_- \rangle$$

Density Matrix

Earlier we showed that:

$$\langle A(t) \rangle = \sum_{n,m} a_n^*(t) a_m(t) \langle n|A|m \rangle = \text{Tr}[A\rho(t)]$$

We also showed that

$$\rho(t) = U\rho(0)U^\dagger$$

$$\langle A(t) \rangle = \text{Tr}[A(t)\rho(0)].$$

ρ_0 (or $\rho(0)$ or $\rho(-\infty)$ or ρ_{eq}) is the equilibrium canonical density matrix.

$$\rho_0 = \frac{e^{-\beta H}}{Z} \quad Z = \text{Tr}(e^{-\beta H})$$

for $H|n\rangle = E_n|n\rangle$

$$Z = \sum_n \langle n|e^{-\beta H}|n\rangle = \sum_n e^{-\beta E_n}$$

$$\begin{aligned} (\rho_0)_{nm} &= \frac{1}{Z} \langle n|e^{-\beta H}|m\rangle = \frac{1}{Z} e^{-\beta E_n} \delta_{n,m} \\ &= p_n \delta_{n,m} \end{aligned}$$

So, the ensemble averaged expectation value (at equilibrium) is

$$\langle A(t) \rangle = \text{Tr}[A(t)\rho_0]$$

or for correlation functions:

$$\begin{aligned} C_{AA} &= \sum_n p_n \langle n|A(t)A(0)|n\rangle & p_n &= \langle n|\rho_0|n\rangle \\ &= \text{Tr}(\rho_0 A(t)A(0)) \\ &= \text{Tr}(A(t)A(0)\rho_0) \end{aligned}$$

More on Stationary Processes¹

We stated that correlation functions are stationary: they do not depend on the absolute point of observation, but rather the time-interval between observations.

Let's look at this a bit more: the ensemble average value of A can be expressed as a time-average or an ensemble average. For an equilibrium system:

$$\bar{A} = \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-T}^T dt A_i(t) \quad \bar{\dots} = \text{time - average}$$

$$\langle A \rangle = \sum_n \frac{e^{-\beta E_n}}{Z} \langle n | A | n \rangle \quad \langle \dots \rangle = \text{equil. ensemble average}$$

These quantities are equal for an ergodic system $\langle A \rangle = \bar{A}$. We assume this property for our correlation functions.

So, the correlation of fluctuations can be written:

$$\overline{A(t)A(0)} = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T d\tau A_i(t+\tau)A_i(\tau)$$

$$\langle A(t)A(0) \rangle = \sum_n \frac{e^{-\beta E_n}}{Z} \langle n | A(t)A(0) | n \rangle$$

¹ See McQuarrie, p. 553