

I. Density fluctuations

- 1) Derive the equilibrium value of density fluctuations

$$S(k) = \frac{1}{N} \langle |\delta\rho_k|^2 \rangle = 1 + \rho_0 \hat{h}(k)$$

- 2)
- $S(k=0) = \rho_0 k_B T \chi_T$
- , where

$$\chi_T = \frac{1}{\rho} \left(\frac{\partial \rho}{\partial P} \right)_T$$

is the isothermal compressibility.

- 3) Show
- $2 \operatorname{Re}[\hat{F}(k, z = i\omega)] = S(k, \omega)$
- where
- $\hat{F}(k, z) = \int_0^\infty e^{-zt} F(k, t) dt$
- is the Laplace transform.

- 4) To first order in
- k
- , we obtain

$$F(k, t) = S(k) \left[\left(1 - \frac{1}{\gamma}\right) e^{-ak^2 t/\gamma} + \frac{1}{\gamma} \cos(kc_s t) e^{-k^2 \Gamma t} \right].$$

Use the expression for $F(k, t)$ to derive $S(k, \omega)$.

- 5) Make a plot of
- $S(k, \omega)$
- for a typical value of
- k
- . (Ref: Reichl p 544-562)

II. The transverse current is defined as

$$J_k = \sum_j^N v_{jx}(t) e^{-ikz_j(t)}.$$

- 1) Show
- $\langle |J_k|^2 \rangle = \frac{Nk_B T}{M}$
- (note
- $\langle v_{ix} v_{jx} \rangle = \delta_{ij} \langle v_x^2 \rangle$
-)

- 2) Use the Navier-Stokes equation to show
- $J_k(t) = J_k(0) e^{-\nu k^2 t}$
- where
- $\nu = \frac{\eta}{\rho_0 m}$
- .

- 3) Prove

$$\sum_{ij} \langle v_{ix}(t) v_{jx}(0) [z_i(t) - z_j(0)]^2 \rangle = \langle |A(t) - A(0)|^2 \rangle \frac{1}{m^2}$$

when $A(t) = m \sum_i v_{ix}(t) z_i(t)$.

- 4) *Derive

$$\eta = \frac{1}{V k_B T} \lim_{t \rightarrow \infty} \frac{1}{2t} \langle |A(t) - A(0)|^2 \rangle.$$

- 5) Derive

$$\eta = \frac{1}{V k_B T} \int_0^\infty \langle \sigma_{xz}(t) \sigma_{xz}(0) \rangle dt,$$

where $\sigma_{xz} = \frac{d}{dt} A$.

- 6) *Show
- $\langle \sigma_{\alpha\beta} \rangle = PV \delta_{\alpha\beta}$
- (Ref. 5.70 notes, McQuarrie p513-525)

III. Using the mean-free path approximation [McQuarrie p. 523], derive expressions for D and η in terms of the average collision time t_c .

IV. Landau-Placzek ratio (ref. Reichl) f_{L-P}

$$f_{L-P} = \frac{\langle \Delta \rho^2 \rangle_{\text{thermal}}}{\langle \Delta \rho^2 \rangle_{\text{mechanical}}} = \frac{\left(\frac{\partial \rho}{\partial S} \right)_P^2 \langle \Delta S^2 \rangle}{\left(\frac{\partial \rho}{\partial P} \right)_S^2 \langle \Delta P^2 \rangle}.$$

1) Explain this expression.

2) Using the thermal fluctuation relation $P \propto \exp \left\{ -\frac{\beta}{2} [\Delta S \Delta T - \Delta P \Delta V] \right\}$ show

$$\frac{\langle \Delta S^2 \rangle}{\langle \Delta P^2 \rangle} = - \frac{\left(\frac{\partial S}{\partial T} \right)_P}{\left(\frac{\partial P}{\partial V} \right)_S}.$$

3) Derive

$$f_{L-P} = \left(\frac{\partial V}{\partial T} \right)_P \left(\frac{\partial P}{\partial S} \right)_\rho.$$

4) *Show $f_{L-P} = \gamma - 1 = \frac{c_P - c_V}{c_V}$ when $\gamma = \frac{c_P}{c_V}$.

V. *Derivations

1) To second order in k , derive

$$F(k, t) = \left(1 - \frac{1}{\gamma} \right) e^{-ak^2 t / \gamma} + \frac{1}{\gamma} e^{-\Gamma k^2 t} [\cos(c_s k t) + d(k) \sin(c_s k t)],$$

$$\text{where } d(k) = \frac{\Gamma + (1 - \frac{1}{\gamma})a}{c_s} k.$$

2) Prove

$$b = \frac{m\beta}{2} \lim_{\omega \rightarrow 0} \lim_{k \rightarrow 0} \left(\frac{\omega}{k} \right)^4 S(k, \omega)$$

$$\text{when } b = \frac{1}{m\rho_0} \left(\eta_B + \frac{4}{3}\eta \right).$$

VI. *Derive the explicit expression for the intensity of scattered light from classical electromagnetic theory. (Ref. Reichl)

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